Anatomy of critical fluctuations in hadronic matter

Michał Marczenko

University of Wrocław, Poland

References:

[1] **MM**, K Redlich, C. Sasaki PRD 107, (2023) 5, 054046

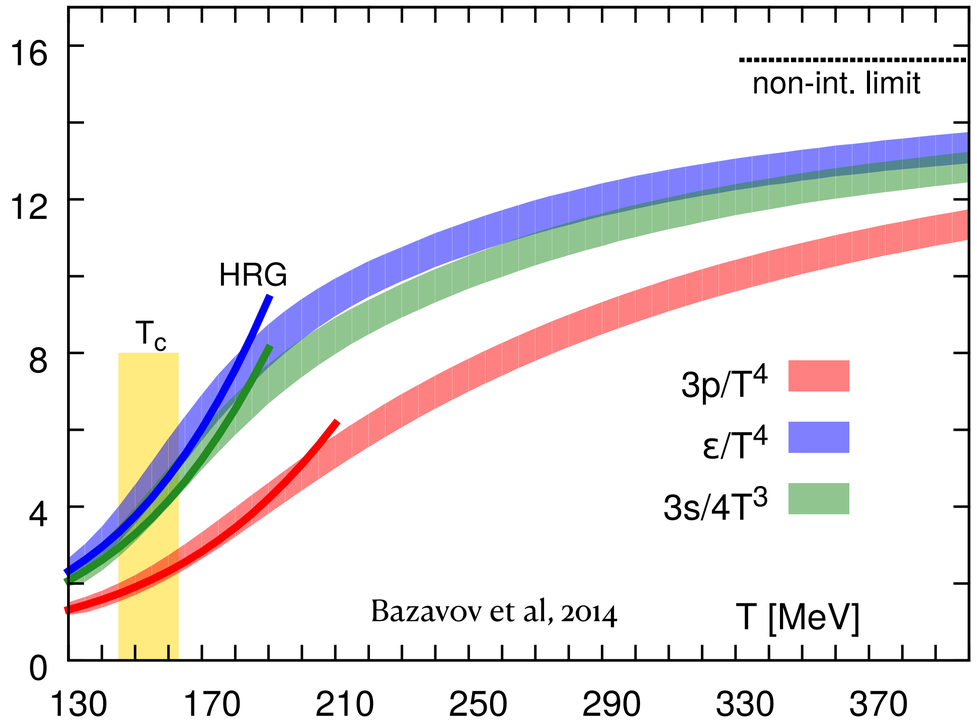
[2] V. Koch, **MM**, K Redlich, C. Sasaki, PRD 109 (2024) 1, 014033

[3] **MM**, PRD 110 (2024) 1, 014018

[4] **MM**, K Redlich, C. Sasaki, PRC (to appear)

65. Jubilee Cracow School of Theoretical Physics, Zakopane, 14-21.06.2025

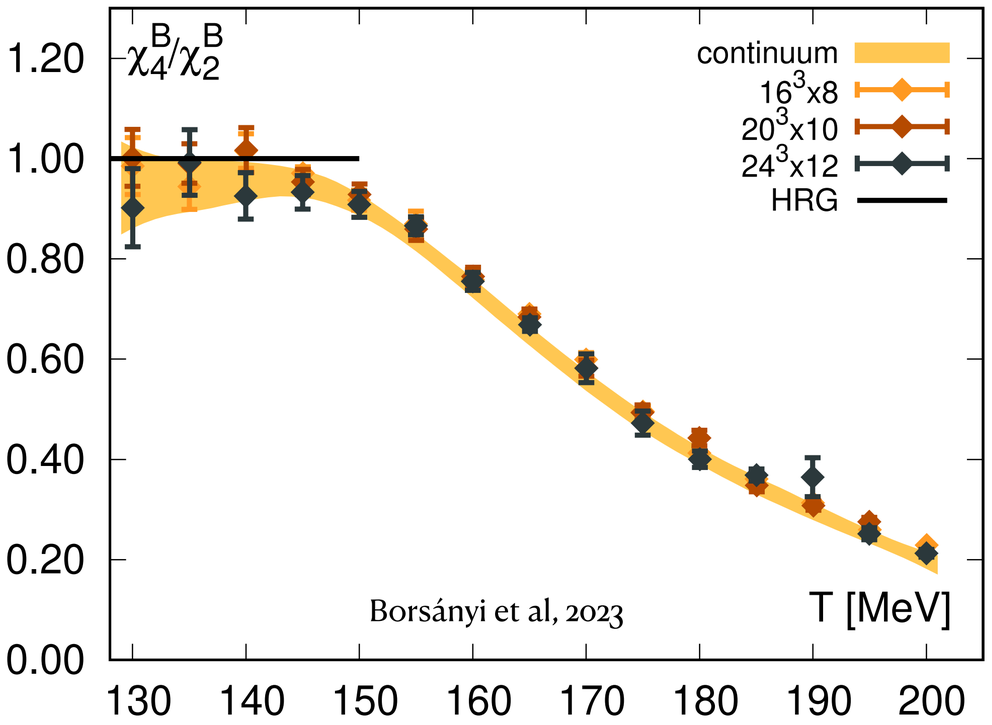
Lattice QCD vs Hadron Resonance Gas



Pressure in the HRG model

$$P^{\mathrm{HRG}} = \sum_{i \in \mathrm{had}} P^{\mathrm{id}}\left(T, \mu_i; m_i\right)$$

Excellent agreement with LQCD EoS up to $\simeq T_c$

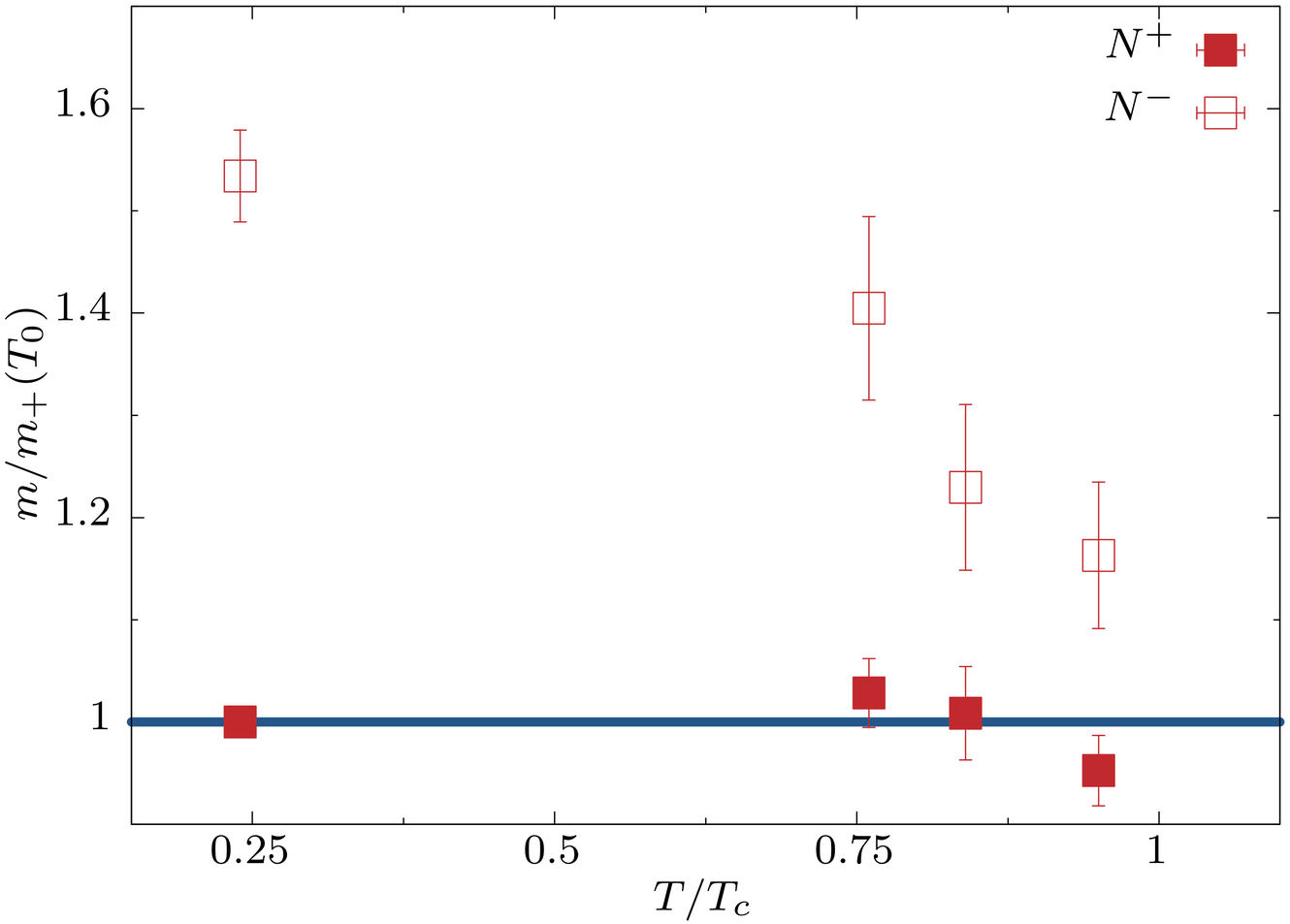


Taylor expansion of LQCD EoS

$$\frac{P}{T^4} = \sum_{k=0}^{\infty} \left(\frac{\mu_B}{T}\right)^k \frac{\chi_k^B}{k!}, \text{ where } \chi_k^B = \frac{\partial^k P/T^4}{\partial\left(\mu_B/T\right)^k}$$

Kurtosis: $\chi_4^B/\chi_2^B \sim B^2$

Parity Doubling in Lattice QCD Aarts et al, 2017, 2019



- N^+ nucleon stays nearly unchanged
- N^- chiral partner drops mass towards T_c
- Chiral partners $N^{\pm}$ degenerate at T_c
- Chiral partners stay massive
- Seen for octet and decouplet of baryons

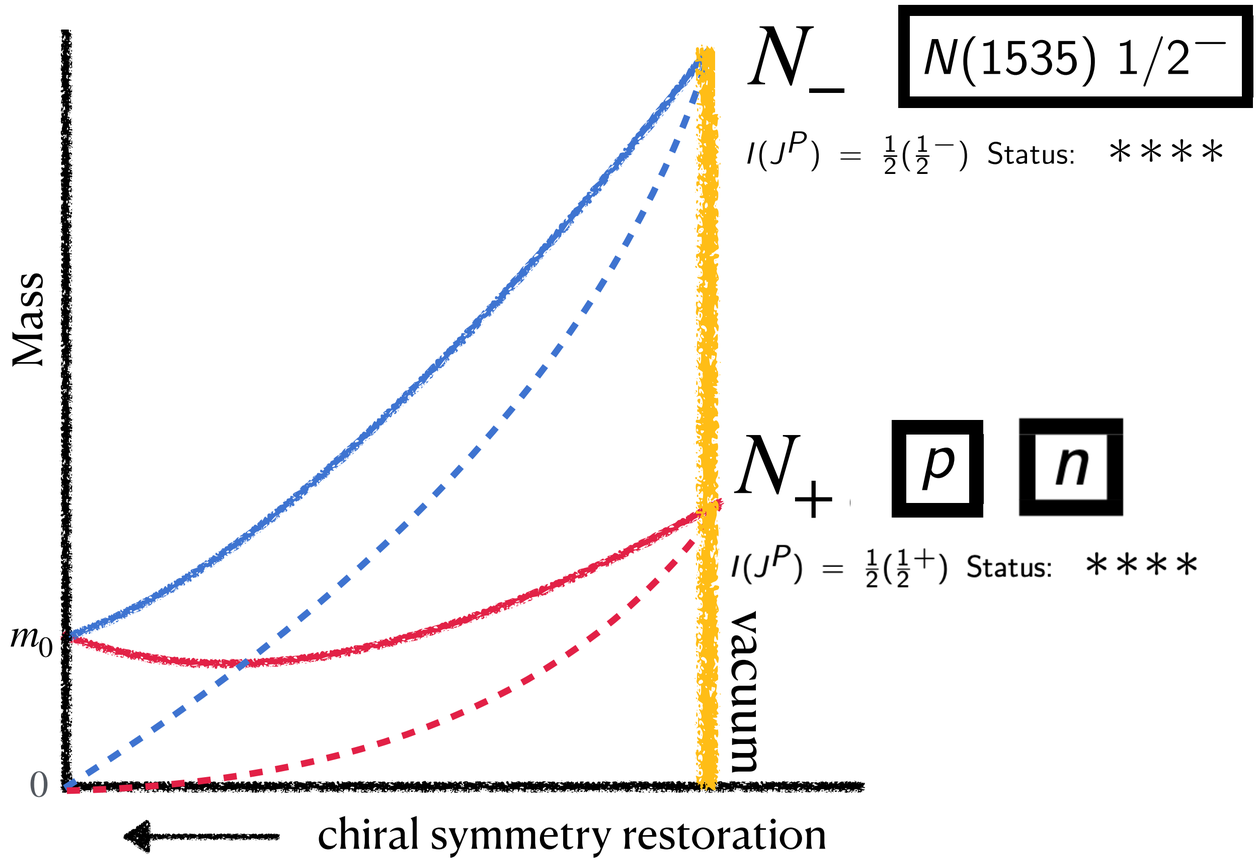
Imprint of chiral symmetry restoration in the baryonic sector

LQCD results still obtained with heavy m_π far from continuum limit

Parity Doublet Model a'la DeTar, Kunihiro 1989

- SU(2) chiral transformation of 2 nucleons → how to assign 2 independent rotation to them?

$$\mathscr{L}_{\text{mass}} \sim m_0 \left(\bar{\psi}_1 \gamma_5 \psi_2 + \bar{\psi}_2 \gamma_5 \psi_1\right) \Longrightarrow M_\pm = \frac{1}{2}\left(\sqrt{4m_0^2 + a^2\sigma^2} \mp b\sigma\right) \xrightarrow{\sigma \to 0} m_0$$

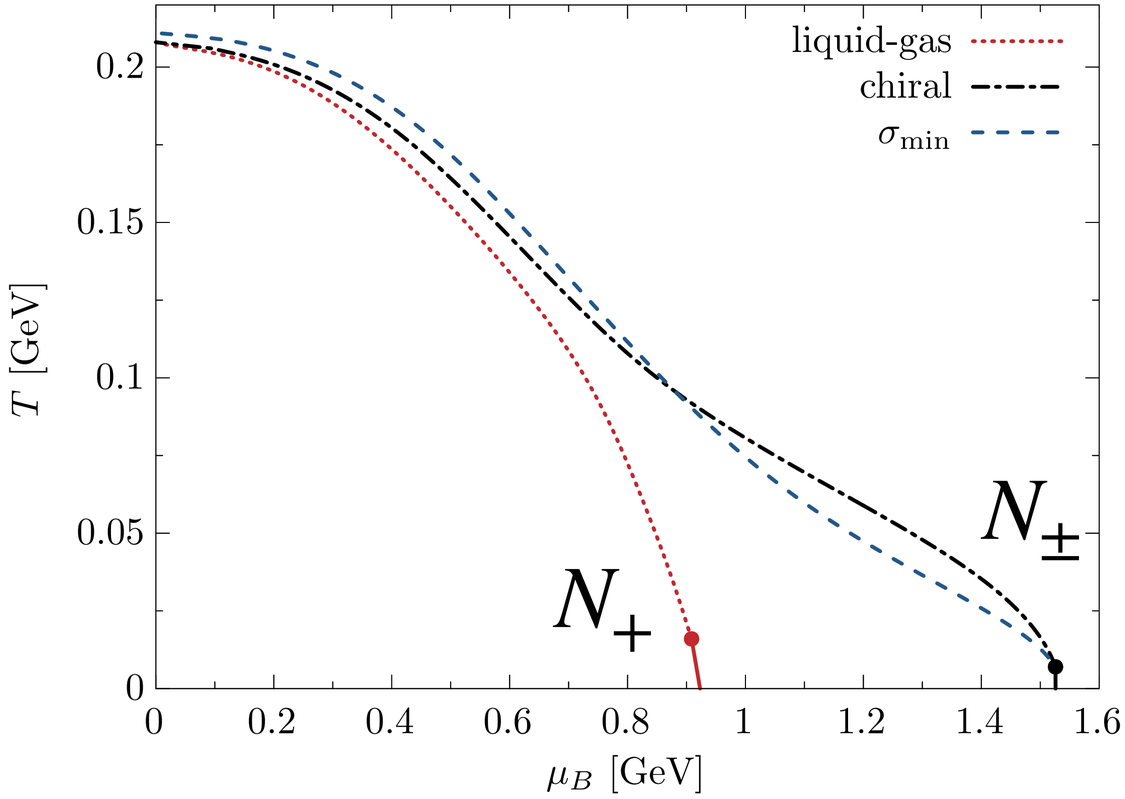


N_- $N(1535)\ 1/2^-$

$I(J^P) = \frac{1}{2}(\frac{1}{2}^-)$ Status: ****

N_+ p n

$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$ Status: ****



For multiplicity $N_B = N_+ + N_-$ Net-baryon number: $\langle N_B \rangle = \langle N_+ \rangle + \langle N_- \rangle$

Second-order fluctuations of the net-baryon number:

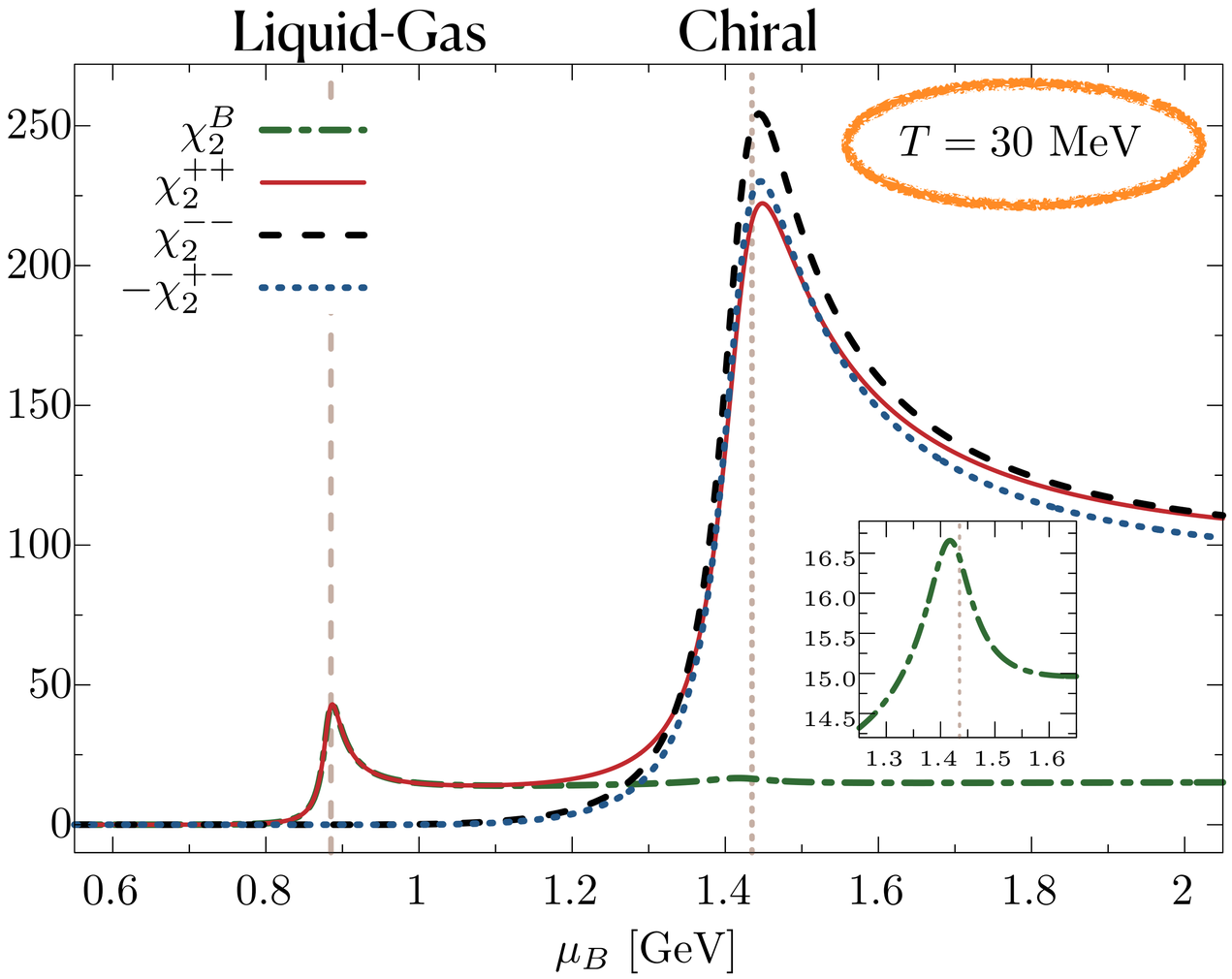
$$\langle \delta N_B \delta N_B \rangle = \langle (\delta N_+)^2 \rangle + \langle (\delta N_-)^2 \rangle + 2\langle \delta N_+ \delta N_- \rangle$$

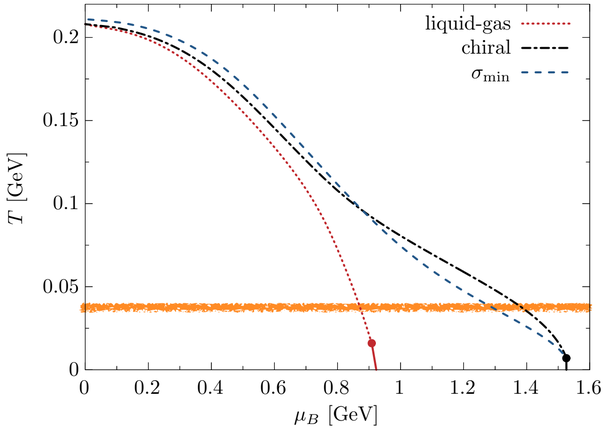
$$\langle \delta N_\alpha \delta N_\beta \rangle = VT^3 \chi_2^{\alpha\beta} \quad \longleftrightarrow \quad \chi_2^{\alpha\beta} = \frac{d^2 P/T^4}{d(\mu_\alpha/T) d(\mu_\beta/T)}$$

$$\chi_2^B = \chi_2^{++} + \chi_2^{--} + 2\chi_2^{+-}$$

- What are the individual contributions of parity partners N_+ and N_-?
- What is the strength and sign of the correlation χ_2^{+-}?
- Is net-proton a good proxy for net-baryon fluctuations? $\chi_2^B = \chi_2^{++} + \chi_2^{--} + 2\chi_2^{+-}$

Fluctuations at liquid-gas and chiral transitions



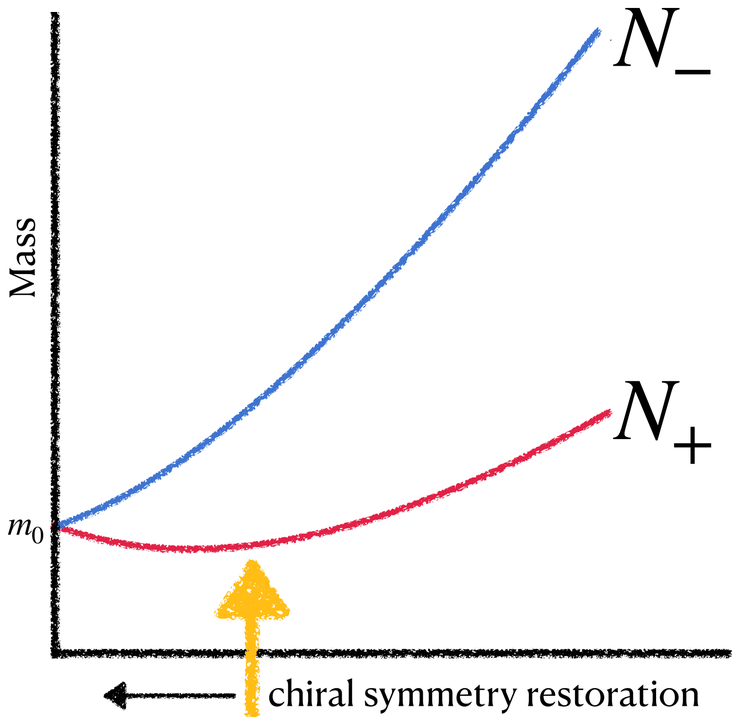


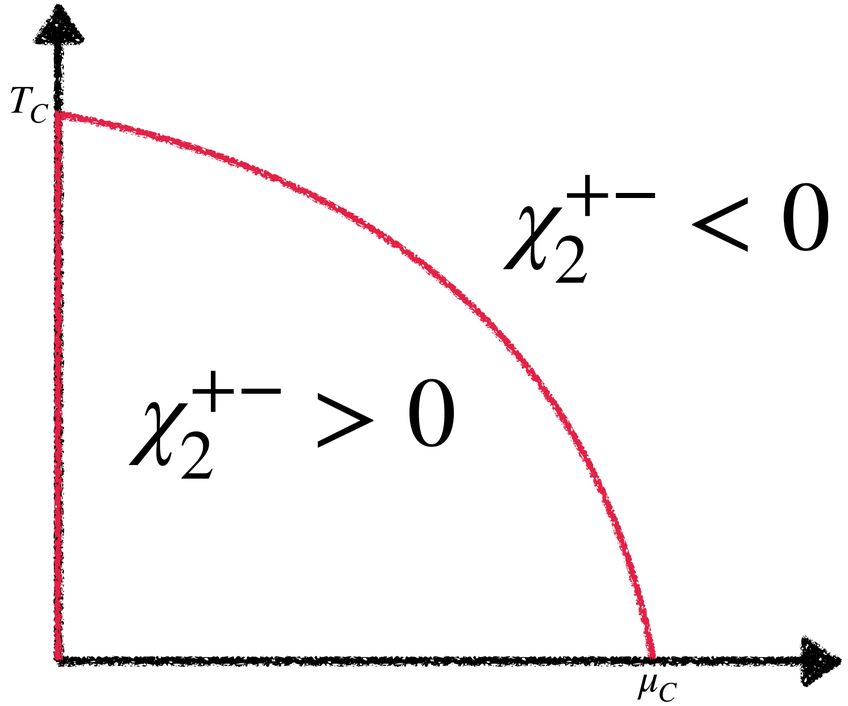
- Liquid-Gas dominated by χ_2^{++}
- Chiral dominated by χ_2^{++} and χ_2^{--}
- Peaks diminished by negative χ_2^{+-}

$$\chi_2^B = \chi_2^{++} + \chi_2^{--} + 2\chi_2^{+-}$$

weak signal in χ_2^B

Idealized behavior of the χ_2^{+-} correlator $\longrightarrow$ no repulsive forces





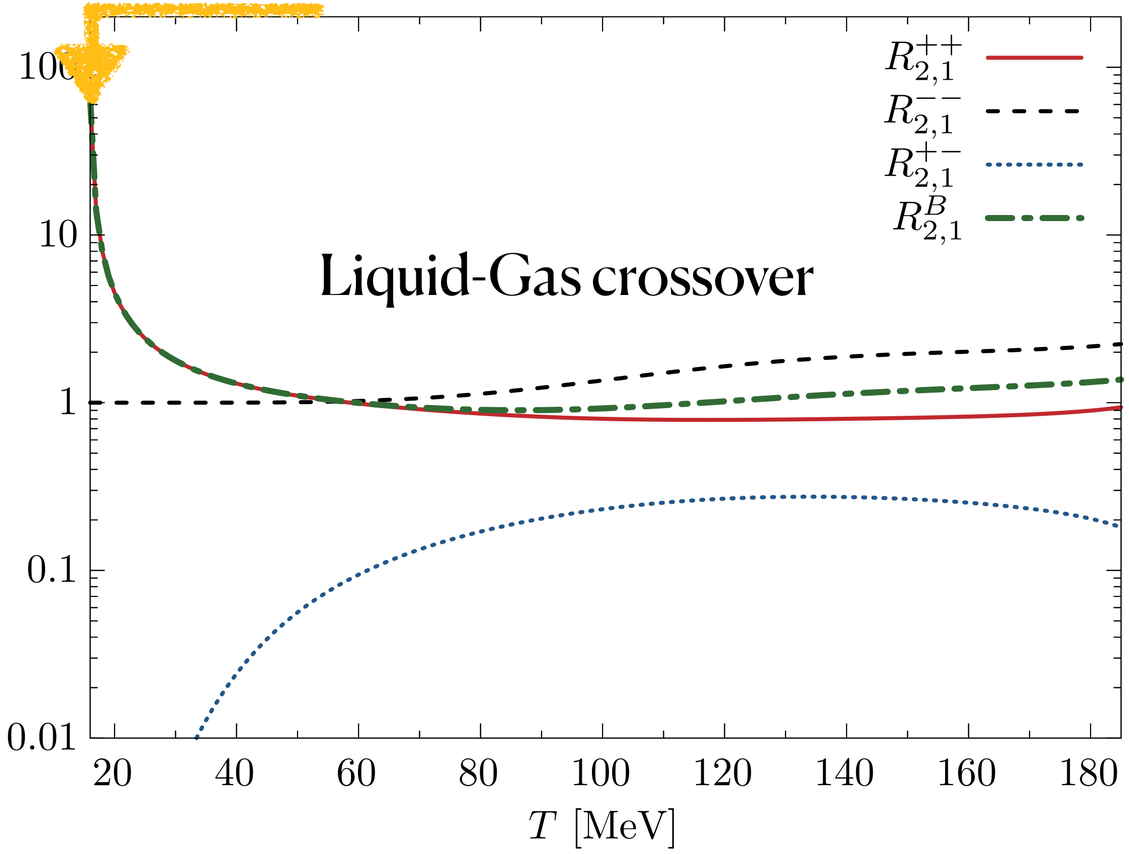
$$\chi_2^{+-} \sim \frac{\partial m_+}{\partial \sigma}\frac{\partial m_-}{\partial \sigma}$$

but also repulsion

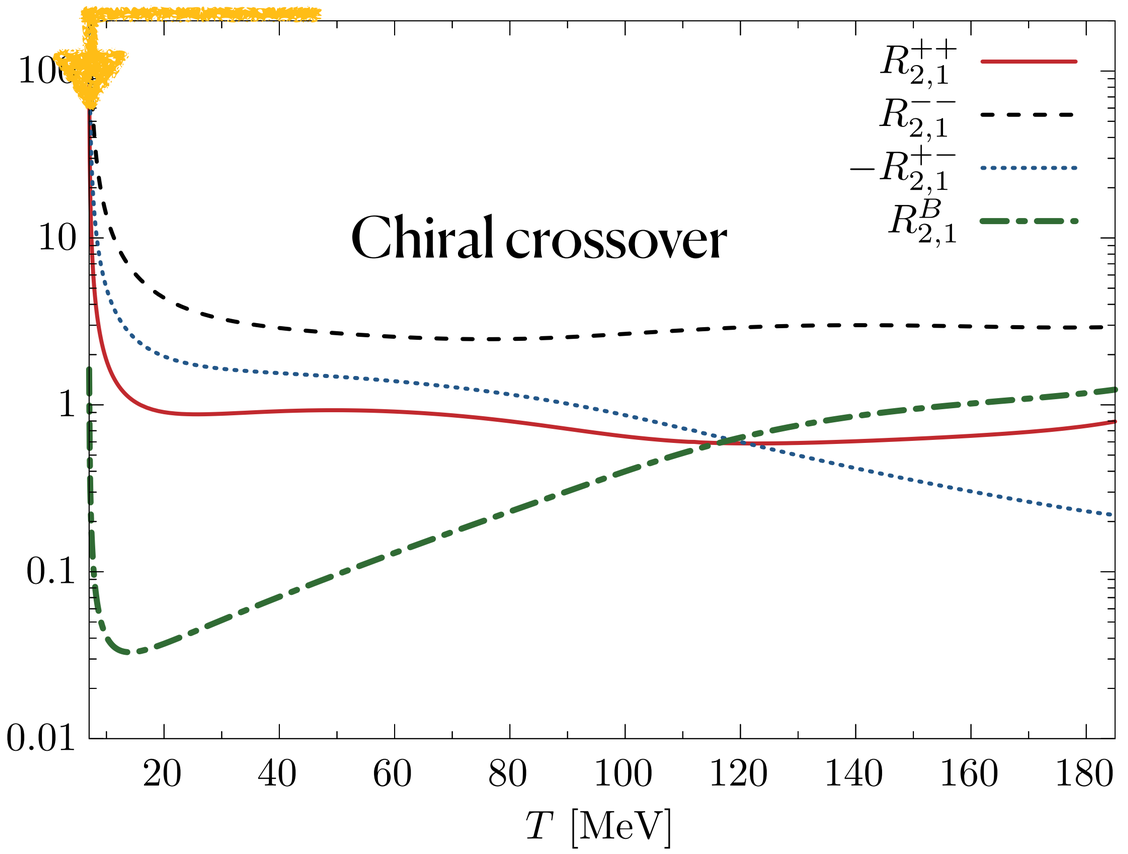
Change of the sign of χ_2^{+-} → chiral phase boundary → interesting quantity to calculate in LQCD

$R_{2,1} = \chi_2/\chi_1$ along phase boundary

Liquid-Gas CP



Chiral CP



Fluctuations dominated by positive parity

↓

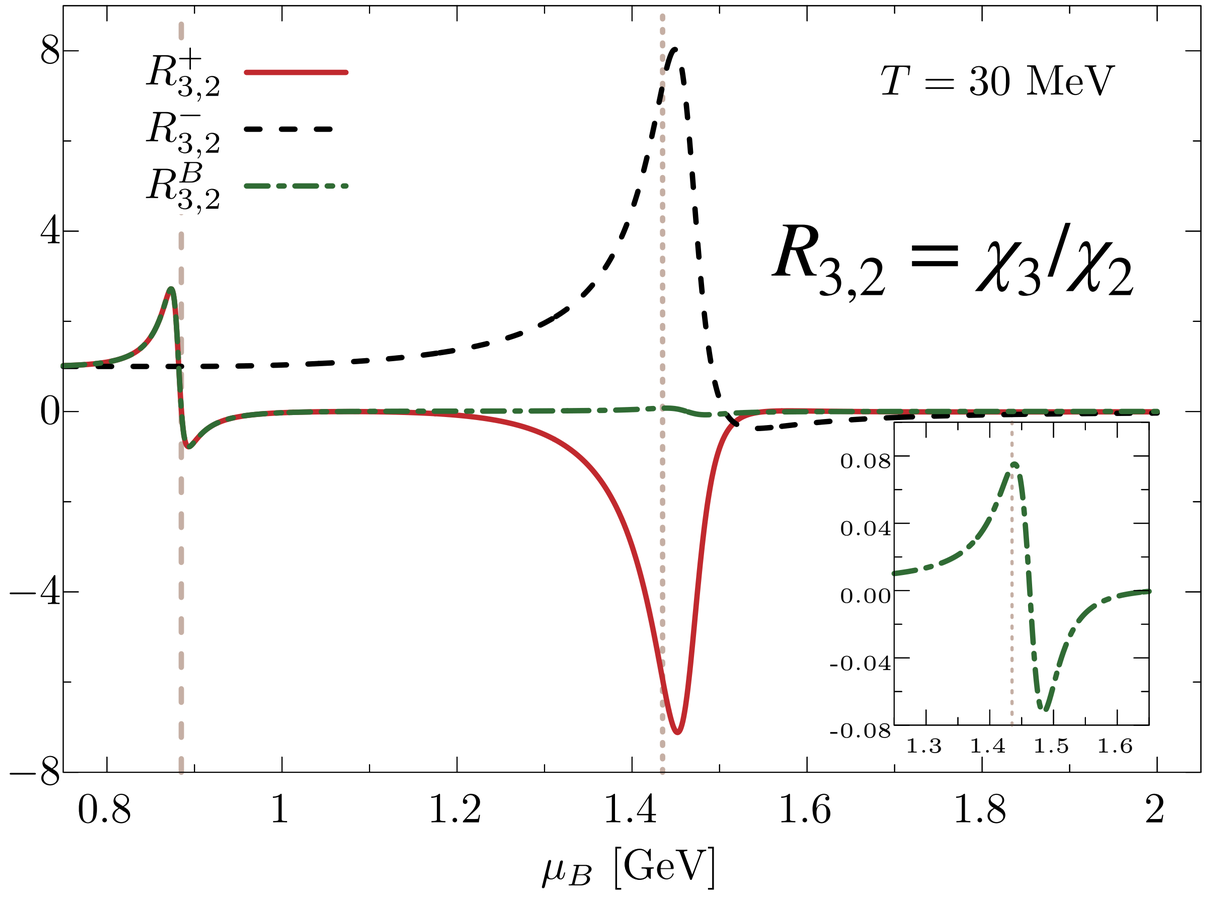
Net-baryon ~ Net-nucleon

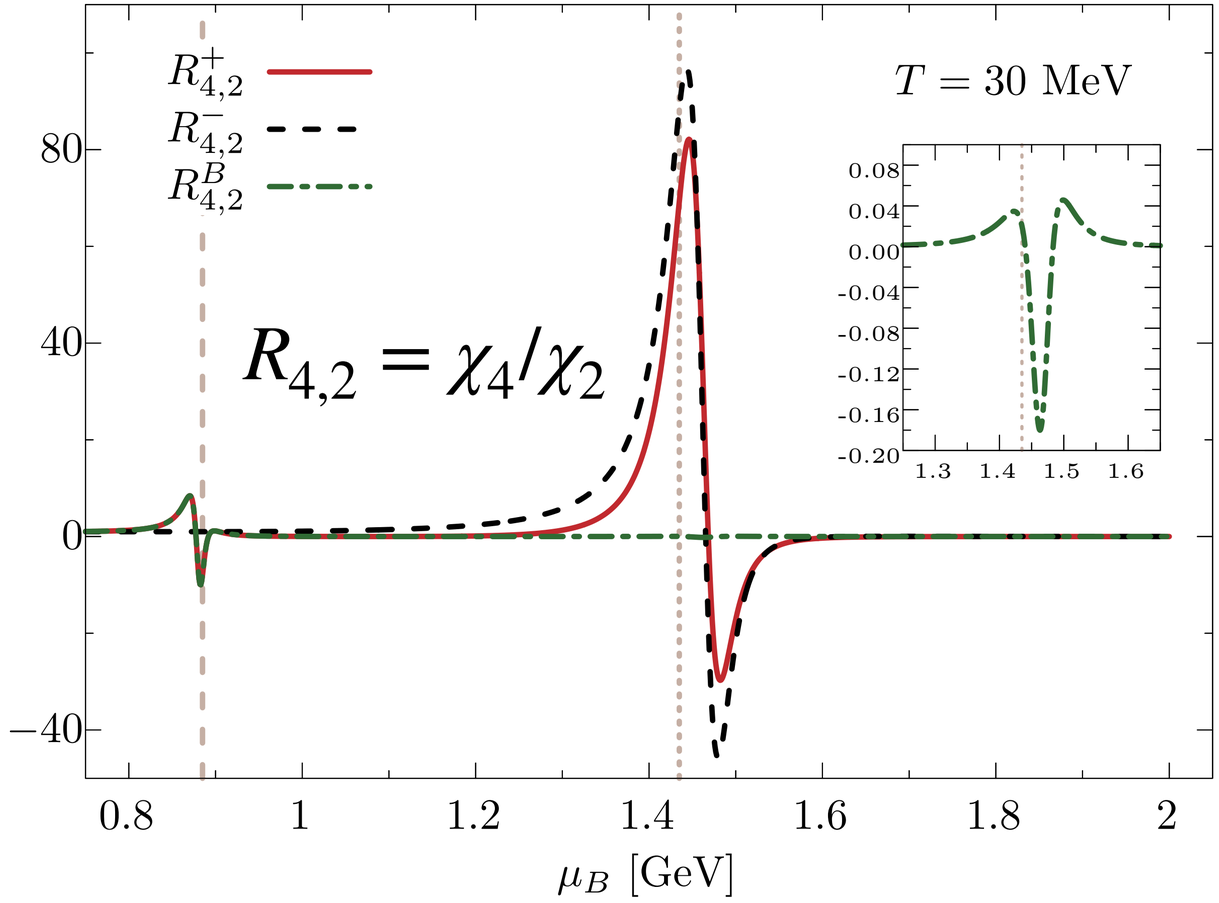
Presence of chiral partners + correlations

↘ ↙

Net-baryon ≪ Net-nucleon

Higher-Order Fluctuations of Parity Partners

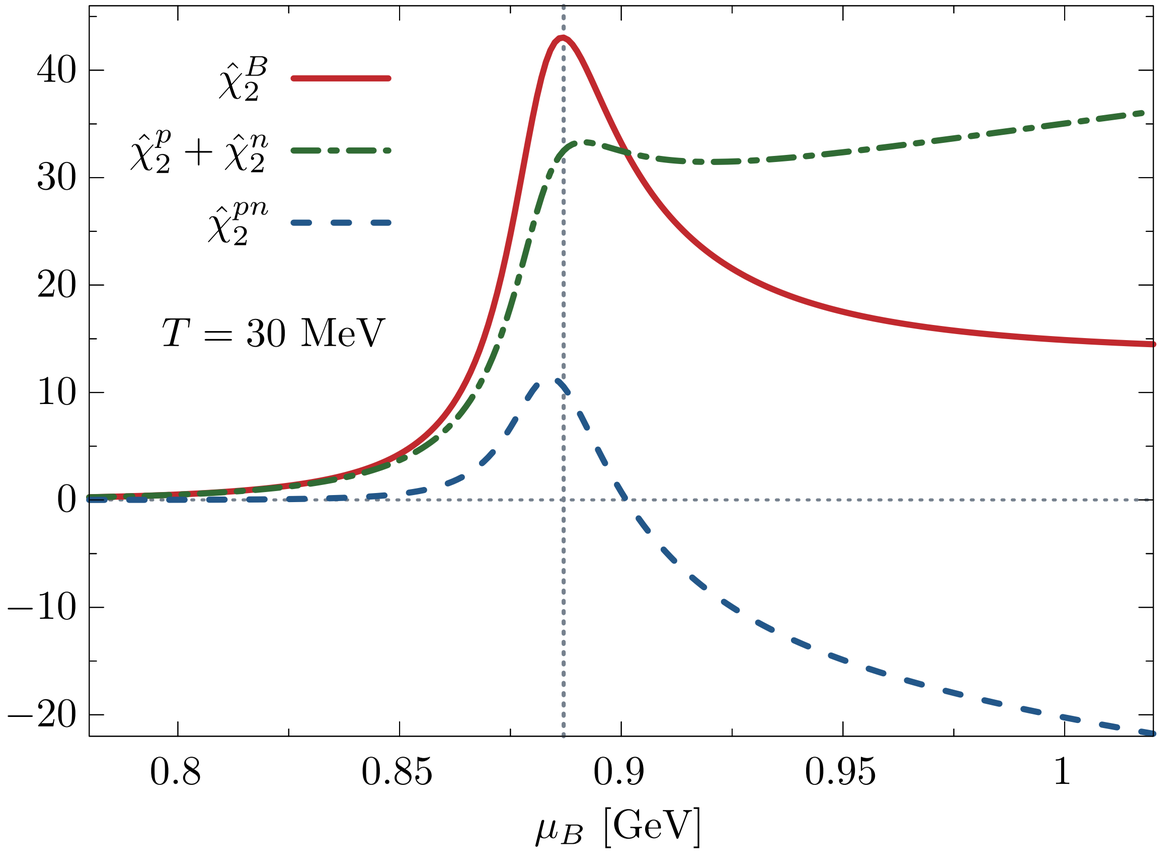


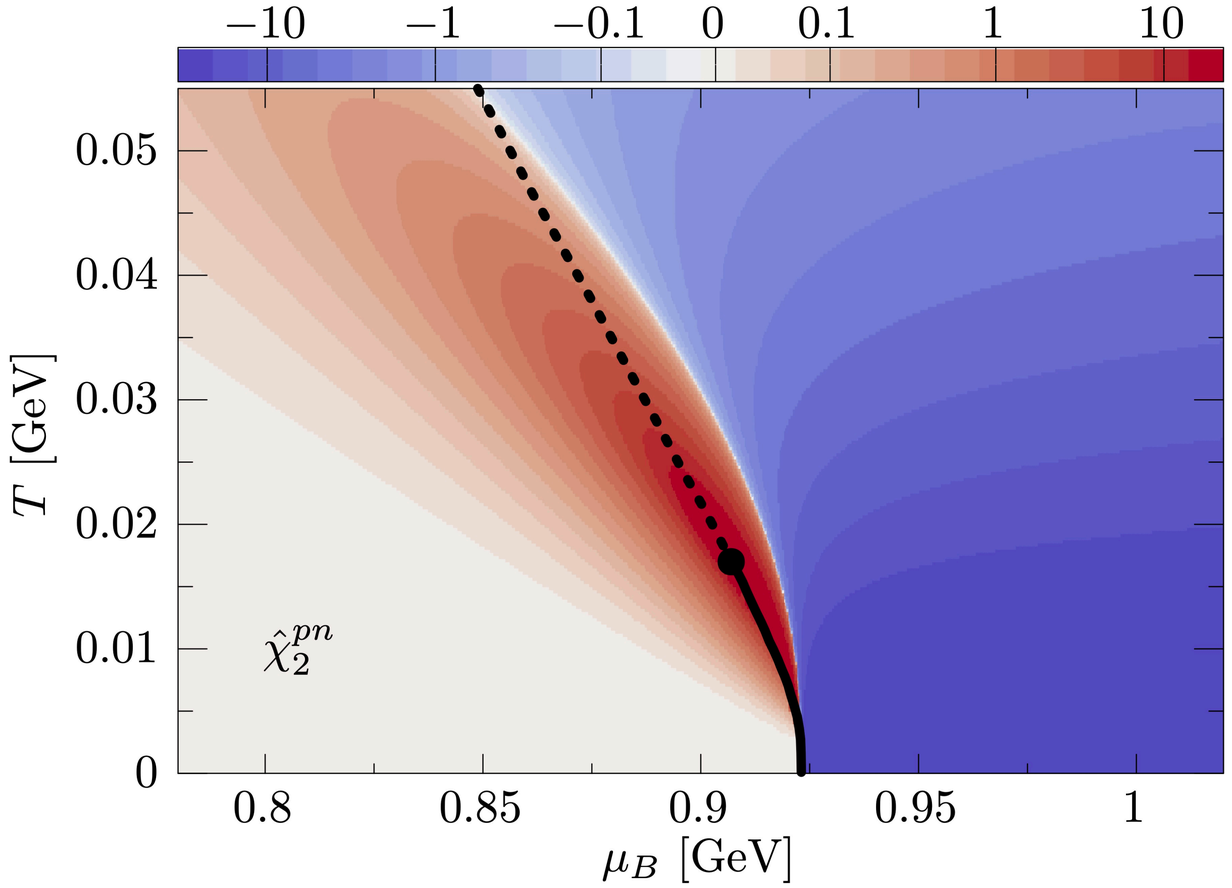


The net-proton fluctuations do not necessarily reflect the net-baryon fluctuations at the chiral phase boundary

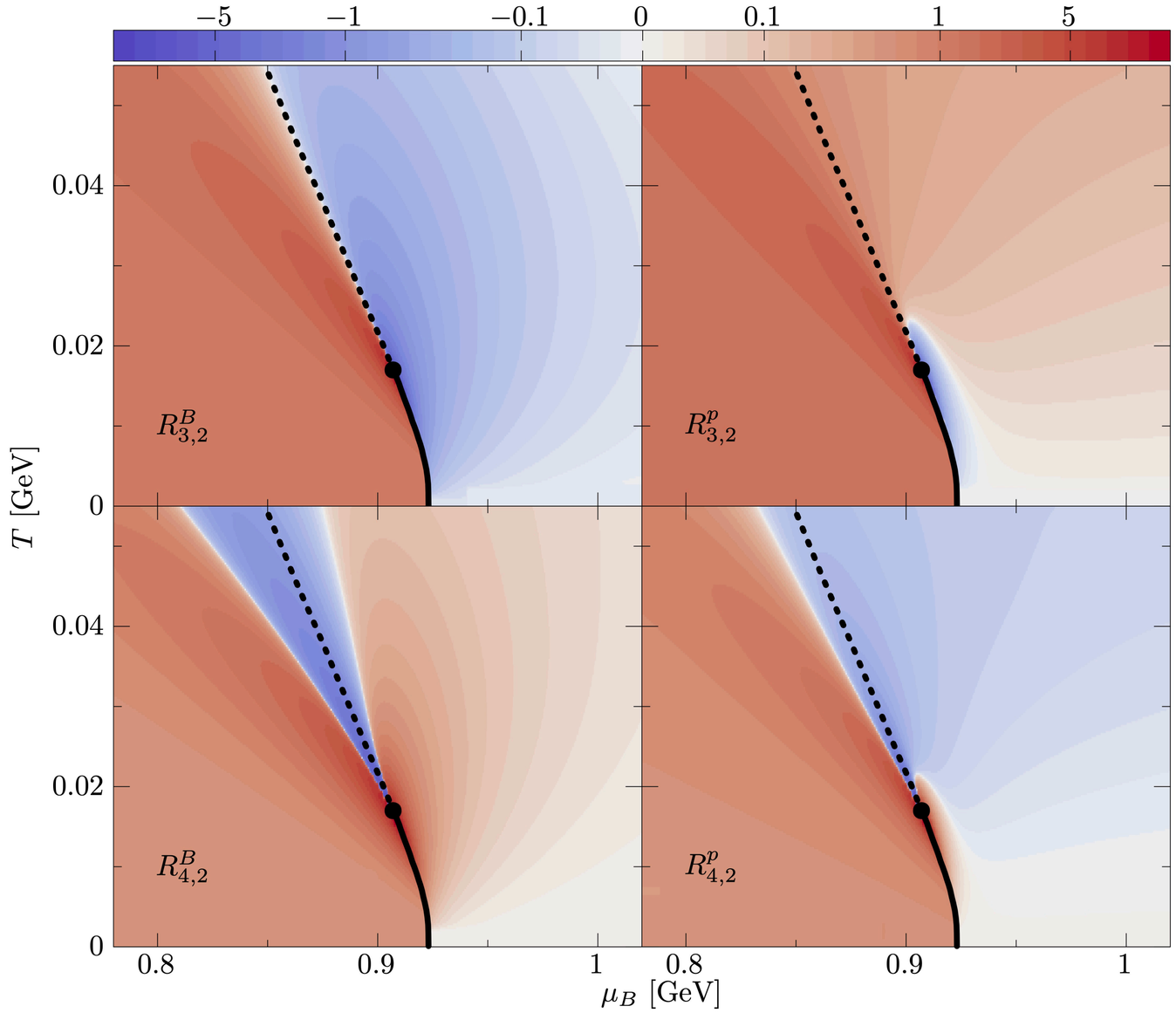
Isospin Correlations Near the Liquid-Gas Transition

$$\chi_2^B = \chi_2^{++} + \ldots \simeq \chi_2^p + \chi_2^n + \chi_2^{pn} \neq 2\chi_2^p$$





$$R^{\alpha}_{n,k} = \frac{\hat{\chi}^{\alpha}_{n}}{\hat{\chi}^{\alpha}_{k}}$$



Differences clearly visible for higher-order fluctuations

Isospin Correlations: Factorial Cumulants

Factorial Cumulants

$$\hat{C}_1^\alpha = \hat{\chi}_1^\alpha$$
$$\hat{C}_2^\alpha = \hat{\chi}_2^\alpha - \hat{\chi}_1^\alpha$$

"Factorial Correlations"

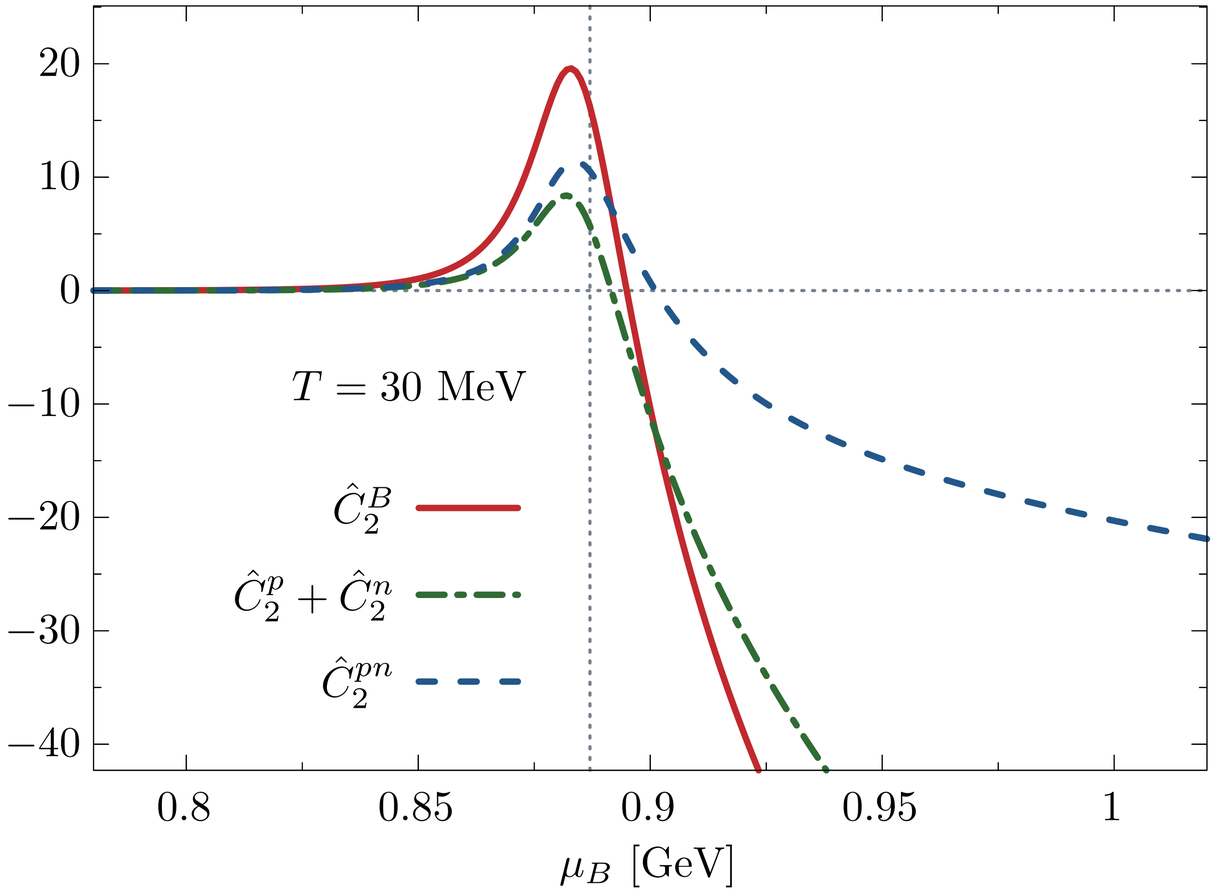
$$\hat{C}_1^{pn} \equiv \hat{C}_1^B - 2\hat{C}_1^p = 0$$

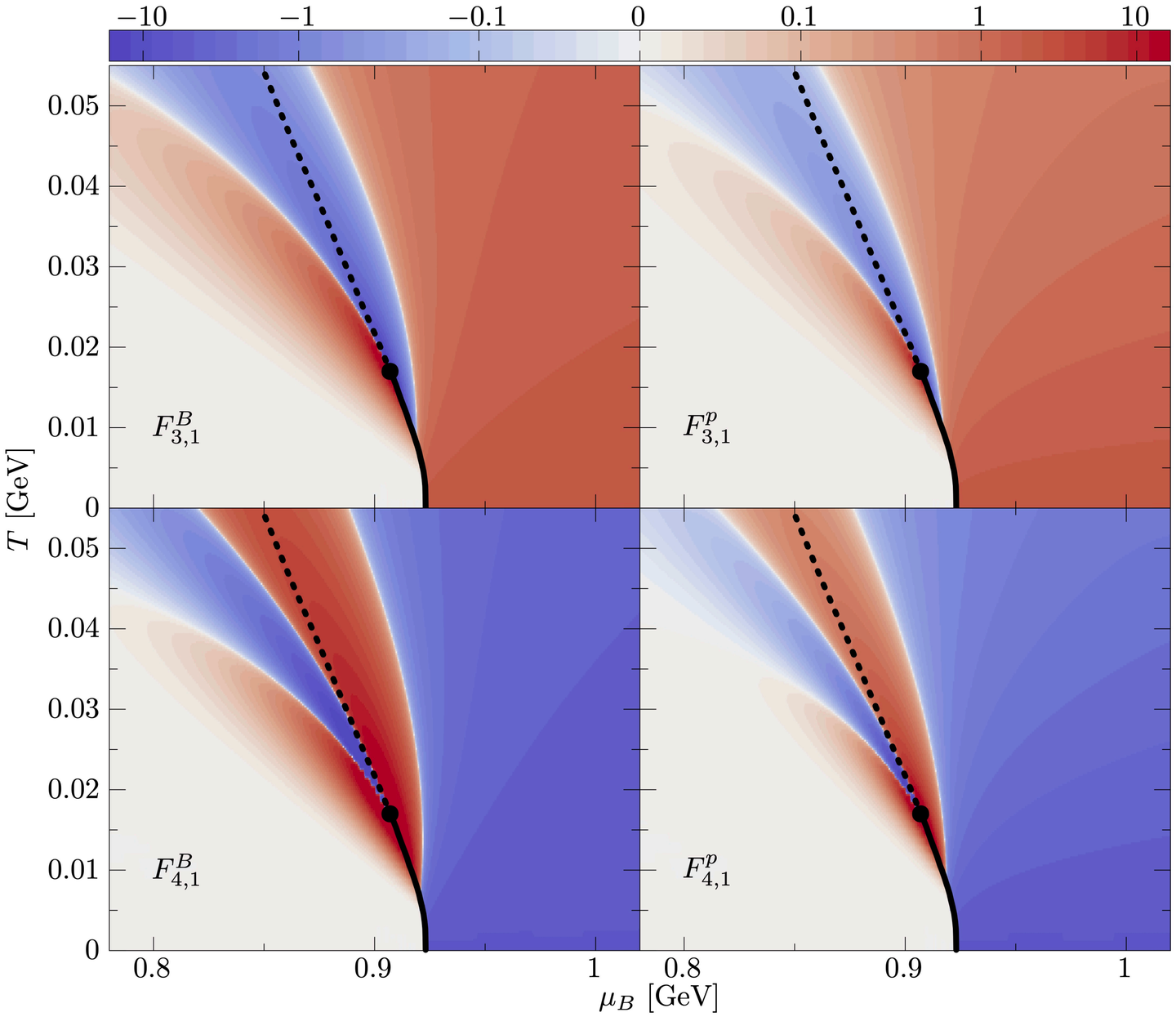
$$\hat{C}_2^{pn} \equiv \hat{C}_2^B - 2\hat{C}_2^p = \hat{\chi}_2^{pn}$$

baryons vs protons:

$$\hat{C}_n^B \sim 2^n \hat{C}_n^p$$ → Same Sign?

Kitazawa, Asakawa (2012)





$$F^{\alpha}_{n,k} = \frac{\hat{C}^{\alpha}_{n}}{\hat{C}^{\alpha}_{k}}$$

In general $\hat{C}^B_n \sim 2^n \hat{C}^p_n$ scaling broken by non-trivial correlations near LG transition

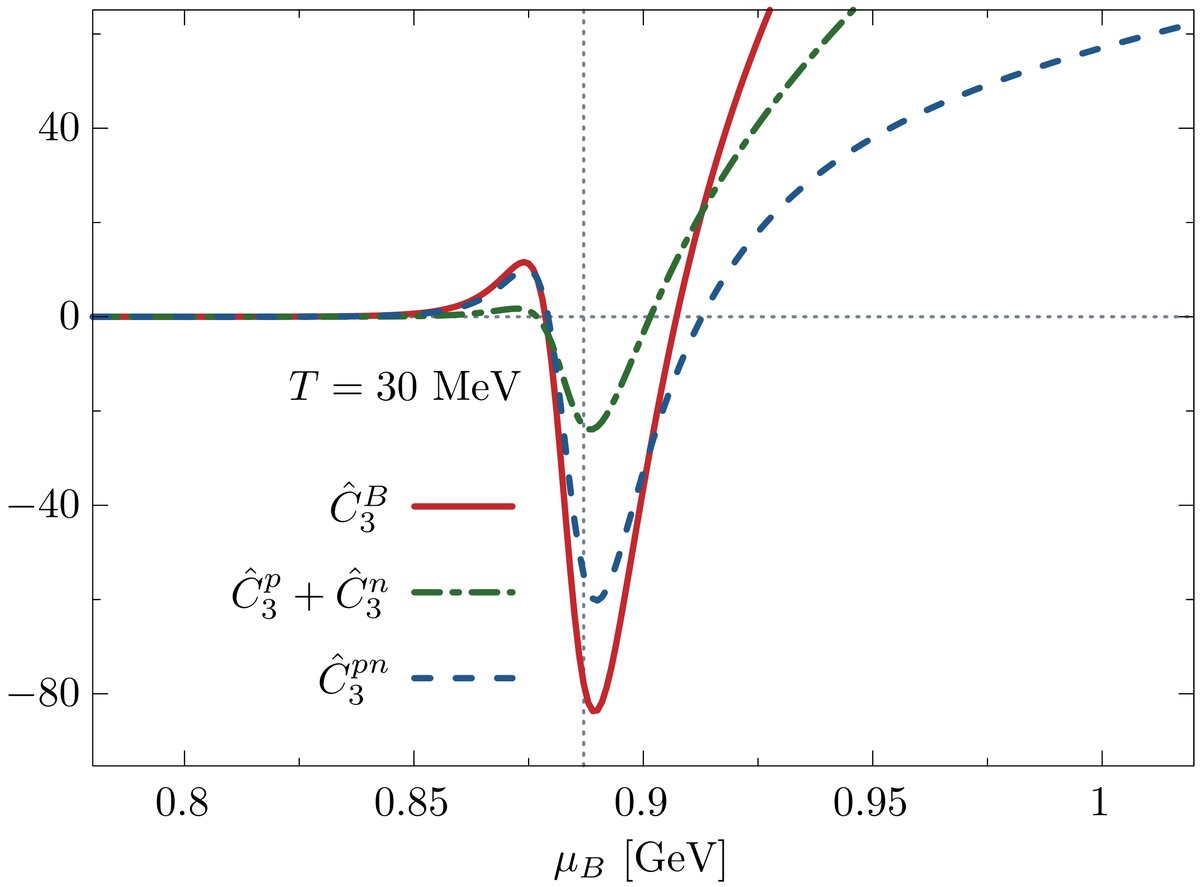
Summary

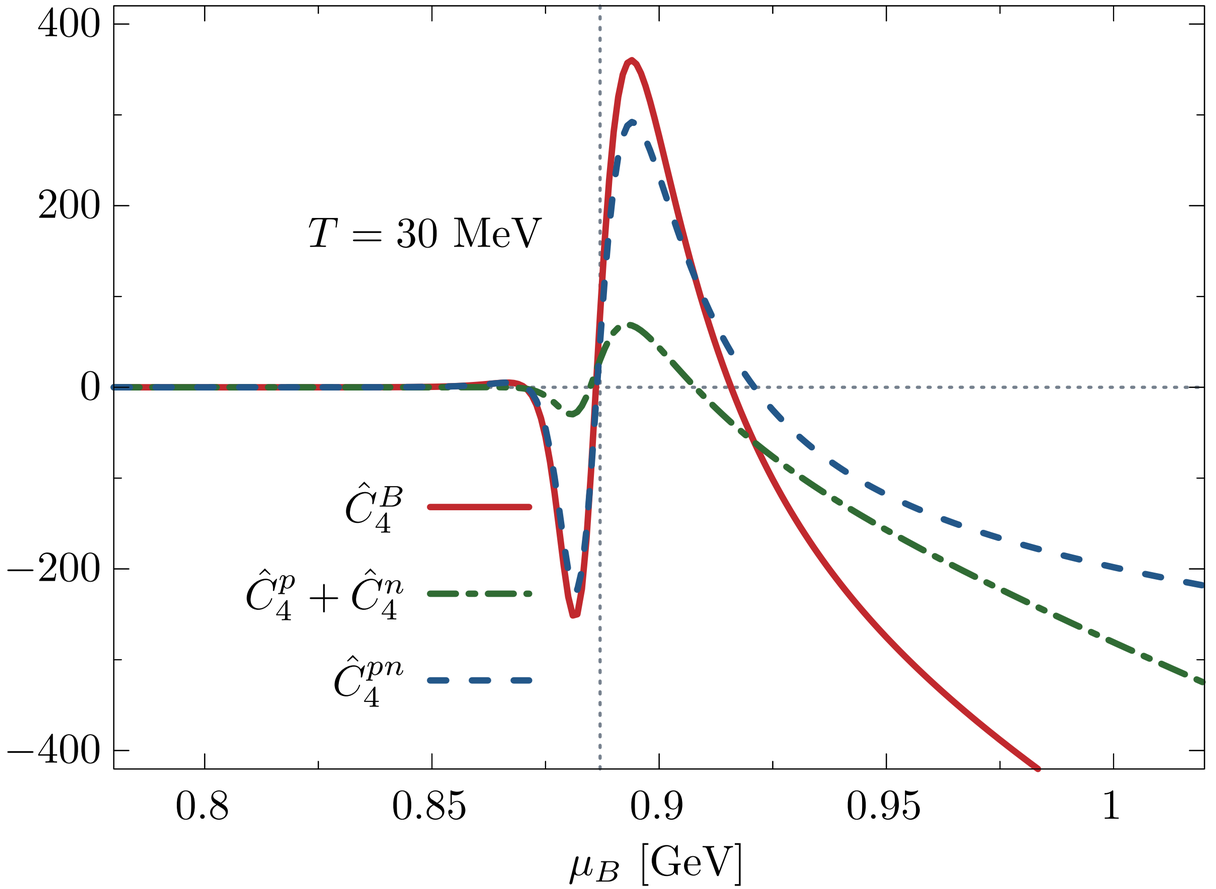
Non-trivial correlations between baryonic chiral partners

Interesting to calculate χ_2^{+-} in other non-perturbative approaches

χ_2^{proton} may not reflect χ_2^B at the chiral or LG phase boundary

Thank You


40
0
−40
−80
$T = 30$ MeV
$\hat{C}_3^B$
$\hat{C}_3^p + \hat{C}_3^n$
$\hat{C}_3^{pn}$
0.8
0.85
0.9
0.95
1
μ_B [GeV]


400
200
0
−200
−400
$T = 30$ MeV
$\hat{C}_4^B$
$\hat{C}_4^p + \hat{C}_4^n$
$\hat{C}_4^{pn}$
0.8
0.85
0.9
0.95
1
μ_B [GeV]

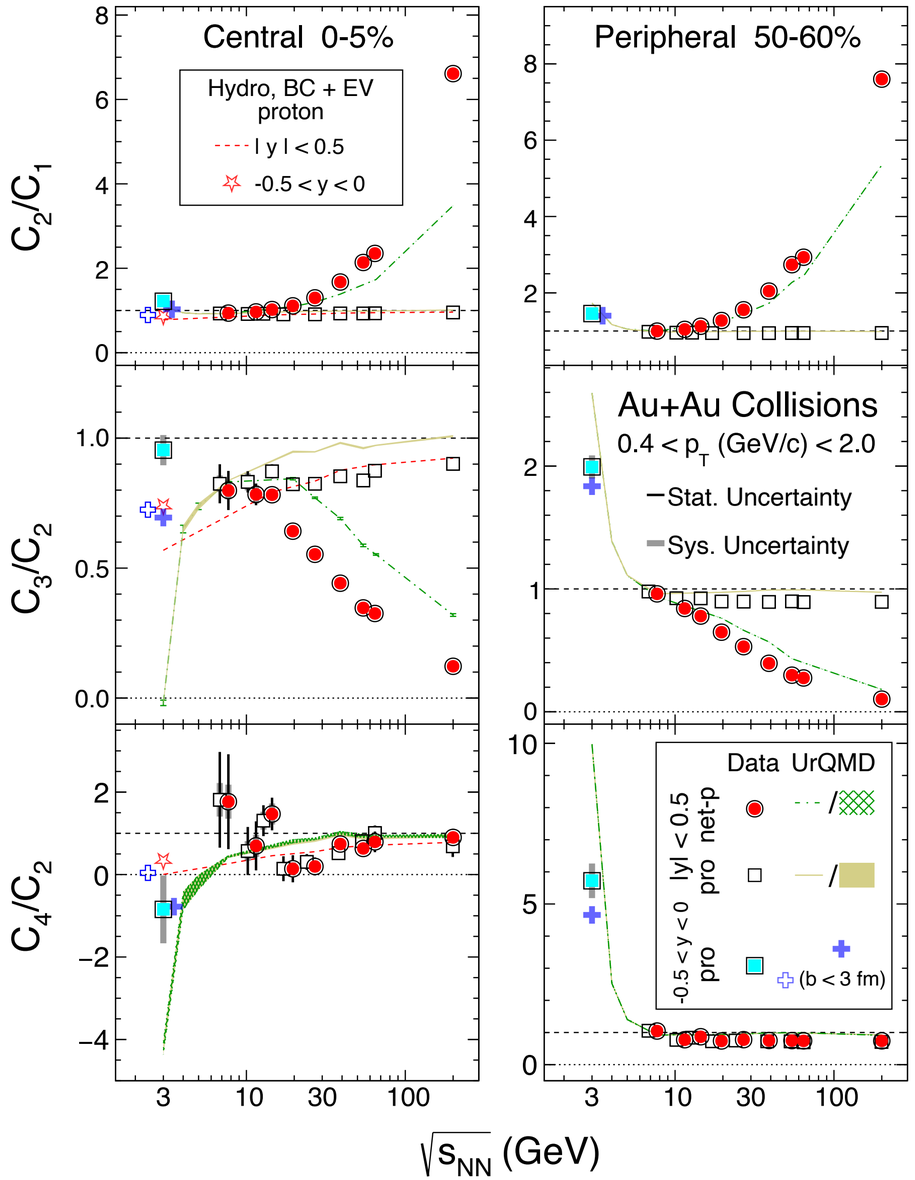
Cumulants vs Susceptibilities

Mean: M	$\langle N_B \rangle$	C_1
Variance: σ^2	$\langle (\delta N_B)^2 \rangle$	C_2
Skewness: S	$\langle (\delta N_B)^3 \rangle / \sigma^3$	$C_3/C_2^{3/2}$
Kurtosis: K	$\langle (\delta N_B)^4 \rangle / \sigma^3 - 3$	C_4/C_2^2

$$C_n \equiv VT^3 \frac{d^n P/T^4}{d\left(\mu_B/T\right)^n}\bigg|_T \longleftrightarrow \chi_n^B \equiv \frac{d^n P/T^4}{d(\mu_B/T)^n}\bigg|_T$$

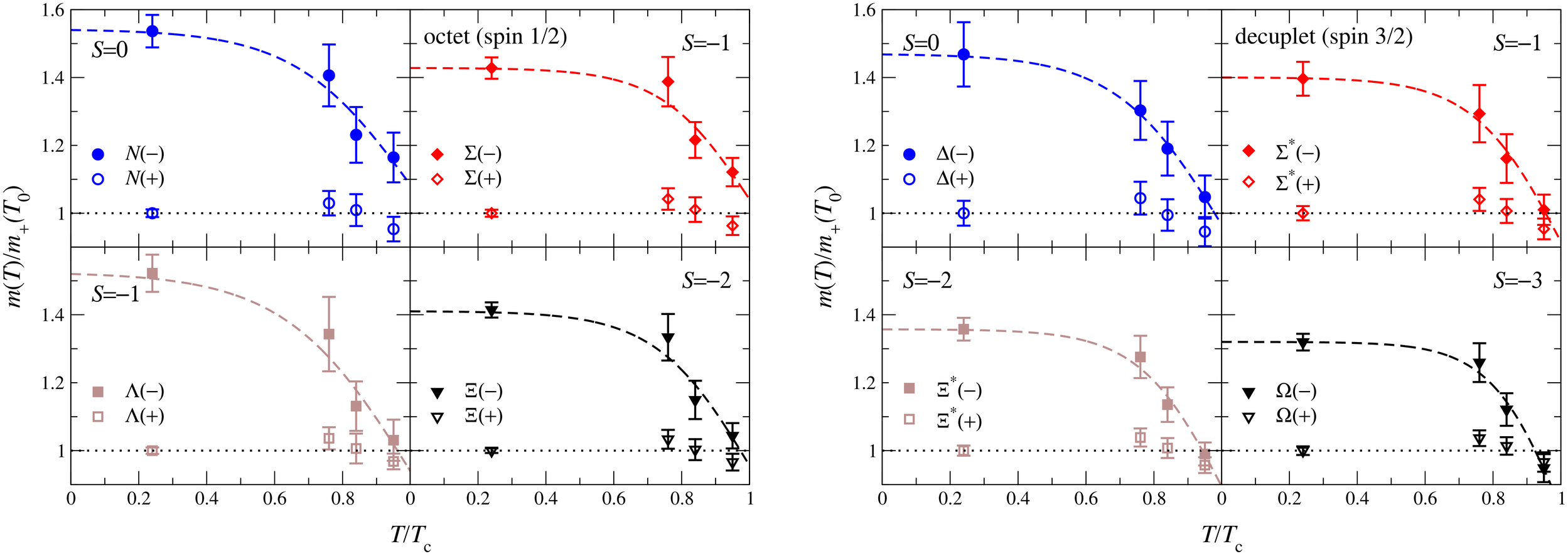
$$C_n = VT^3 \chi_n^B$$

STAR, 2023

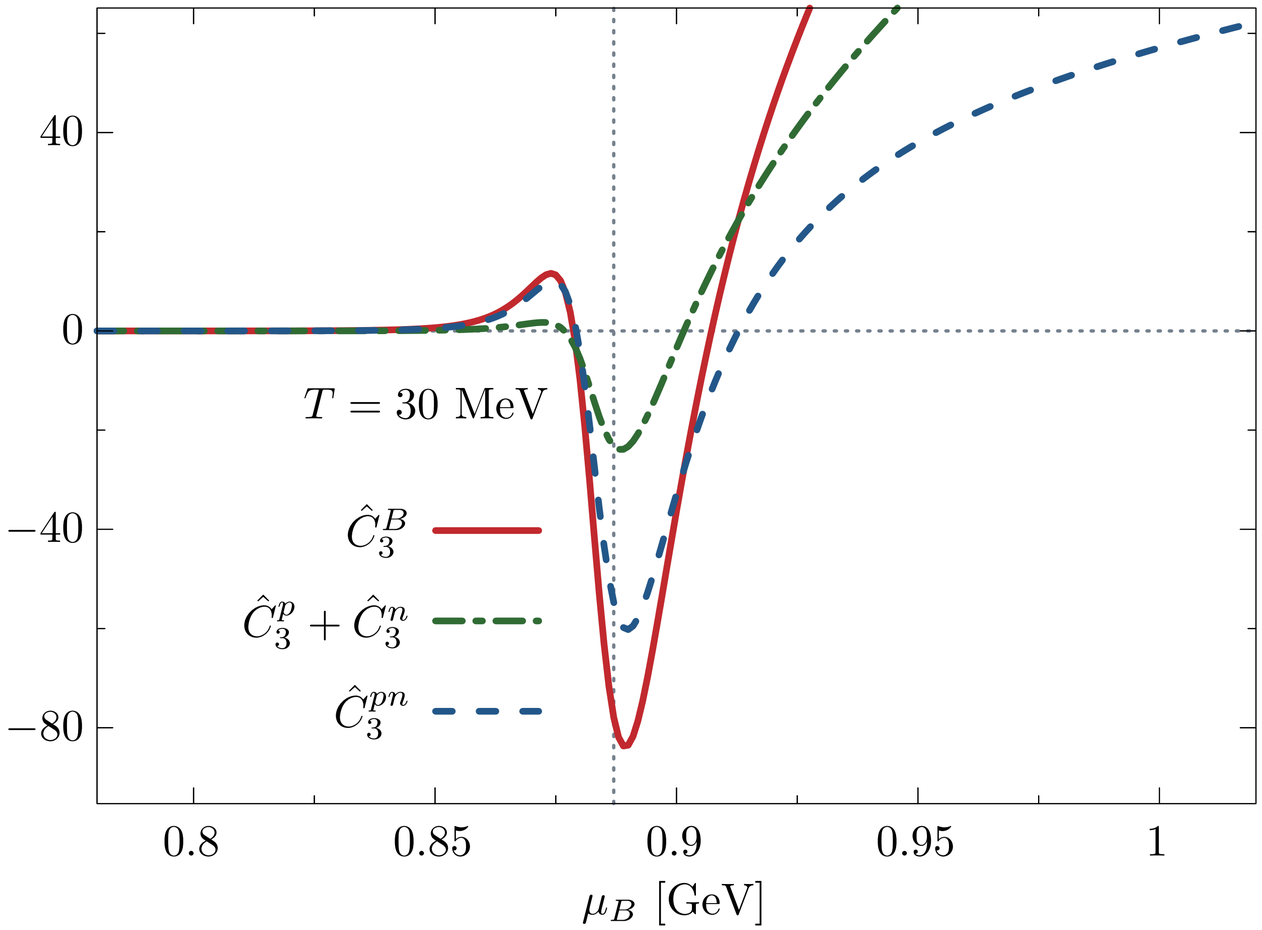


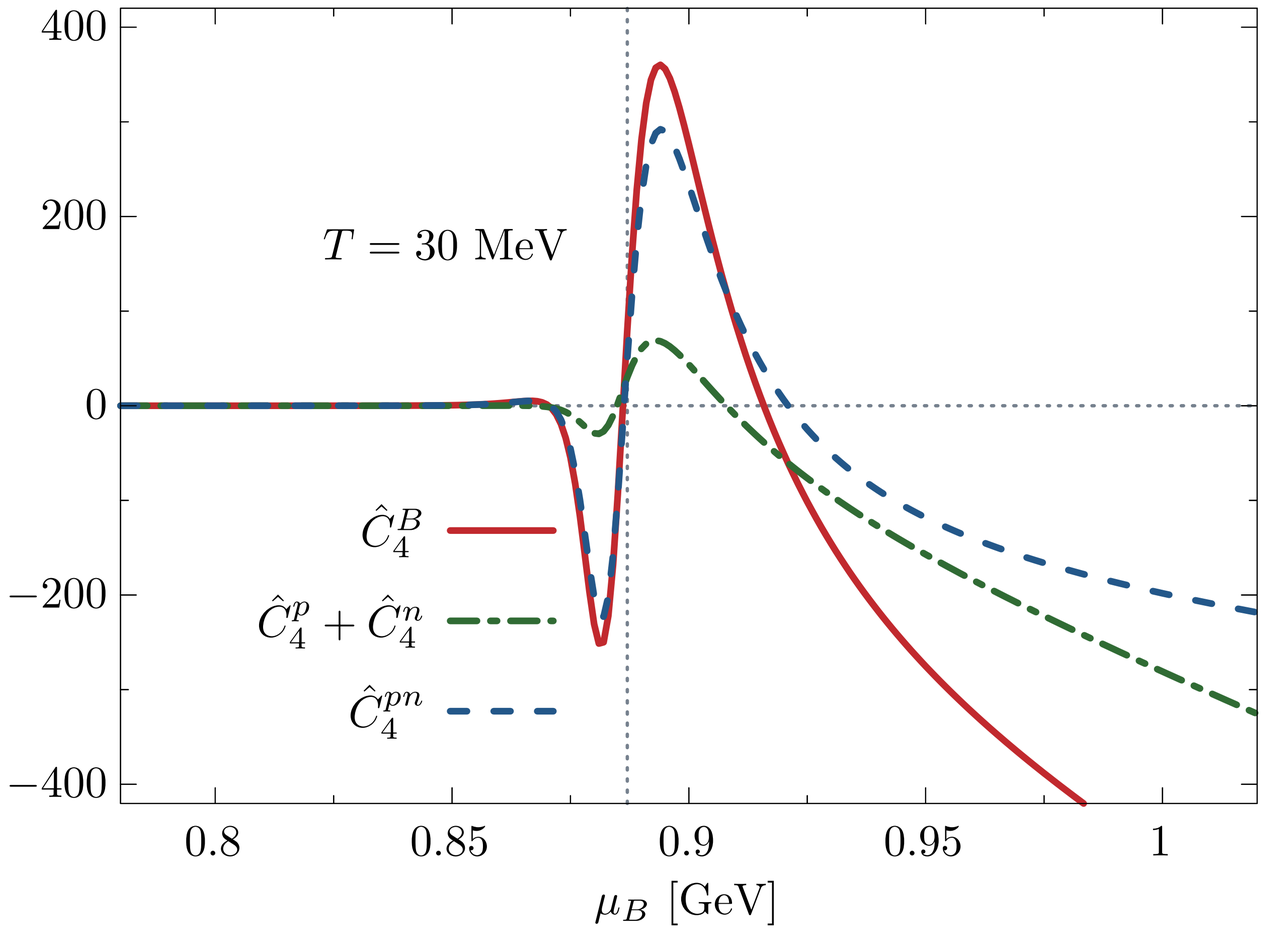
Imprint of chiral symmetry restoration in the baryonic sector

Aarts et al, 2019



Clear evidence for partial restoration of chiral symmetry in the strange baryon sector


40
0
−40
−80
$T = 30$ MeV
$\hat{C}_3^B$
$\hat{C}_3^p + \hat{C}_3^n$
$\hat{C}_3^{pn}$
0.8
0.85
0.9
0.95
1
μ_B [GeV]


400
200
0
−200
−400
$T = 30$ MeV
$\hat{C}_4^B$
$\hat{C}_4^p + \hat{C}_4^n$
$\hat{C}_4^{pn}$
0.8
0.85
0.9
0.95
1
μ_B [GeV]