Gravitational Waves and Stellar-Mass Black-Hole Mergers

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CENTER FOR INTERDISCIPLINARY EXPLORATION AND RESEARCH IN ASTROPHYSICS





 Gain a working knowledge of gravitational wave detections, sources, and implications*

* Primarily relating to stellar-mass compact objects

Overview

01. Theory of Gravitational waves

02. Detecting Gravitational Waves

03. Characterizing GW Sources

04. Catalog of GW Sources

05.

Astrophysical Black-Hole Mergers **06.** Population Inference Agentle introduction to **GWs and binary** sources

$$g_{\alpha\beta} = \eta_{\alpha\beta} + h_{\alpha\beta},$$

Small perturbations *h* to background metric

$$\left(-\frac{\partial^2}{\partial t^2} + \nabla^2\right)\bar{h}^{\alpha\beta} = -16\pi T^{\alpha\beta}.$$

Wave equation w/ strain tensor T (Lorentz gauge)

$$\bar{h}_{lphaeta} = \mathcal{A}\mathbf{e}_{lphaeta} \exp(ik_{\gamma}x^{\gamma}),$$

Wave solutions in vacuum

Schutz and Ricci arxiv:1005.4735

Transverse-Traceless Gauge



Quadrupole formula

$$\bar{h}^{\mathrm{TT}ij} = \frac{2}{r} \stackrel{\cdots}{M}{}^{\mathrm{TT}ij}.$$

At leading order, strain is sourced by time-varying quadrupole moment *M*

$$L_{gw}^{mass} = \frac{1}{5} \left\langle \tilde{M}^{jk} \tilde{M}_{jk} \right\rangle$$

GW luminosity from mass quadrupole

GWs from a quasi-Circular Binary System

$$M_{xx} = \frac{1}{2}\mu R^2 \cos(2\Omega t)$$
$$M^{\text{TT}xx} = M^{xx}/2.$$

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Components of the quadrupole moment



$$\bar{h}^{\text{TT}xx} = -2^{1/3} \frac{\mathcal{M}^{5/3} \Omega_{gw}^{2/3}}{r} \cos \left[\Omega_{gw}(t-r)\right],$$
$$L_{gw} = \frac{4}{5 \cdot 2^{1/3}} \left(\mathcal{M} \Omega_{gw}\right)^{\frac{10}{3}},$$

$$\mathcal{M} := \mu^{3/5} (m_1 + m_2)^{2/5}$$

Solve for strain and luminosity

"Chirp mass"

Energy Balance of GW-Emitting Binary

GW luminosity = orbital energy loss rate

$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left(\left(\frac{5}{96}\right)^3 \pi^{-8} f_{GW}^{-11} \dot{f}_{GW}^3 \right)^{1/5}$$

Straightforward relation between observables and binary parameters!!!!!!!

Chirping binaries are the ultimate astrophysical source

Unimpeded by foregrounds

No calibration to other astro sources

Straightforward astrophysical inference



Effect of spins (lowest order)

$$\chi_{\text{eff}} = \frac{(m_1 \vec{\chi}_1 + m_2 \vec{\chi}_2) \cdot \hat{L}_{\text{N}}}{M}$$



How can we detect these GWs?

Interferometry



Credit: T. Pyle

Interferometer response

$$rac{\delta L(t)}{L} = F_+(heta,\,\phi,\,\psi)h_+(t) + F_ imes(heta,\,\phi,\,\psi)h_ imes(t),$$

$$F_{+} = \frac{1}{2} \left(1 + \cos^{2} \theta \right) \cos 2\phi \cos 2\psi - \cos \theta \sin 2\phi \sin 2\psi,$$

$$F_{\times} = \frac{1}{2} \left(1 + \cos^{2} \theta \right) \cos 2\phi \sin 2\psi + \cos \theta \sin 2\phi \cos 2\psi.$$





Amplitude response of interferometer

(averaged over polarizations)



LIGO Hanford

LIGO Livingston

Operational Planned

Gravitational Wave Observatories

GE0600

KAGRA

Construction for LIGO India

Detectors are noisy!



The amplitude spectral density of noise in detectors

Sources of noise include:

- Shot noise
- Thermal noise
- Ground motion
- Newtonian noise

LVK Phys. Rev. X; 13(4):041039

Statistical distribution of detector noise

$$p(\{n(t_i)\}_i) = \frac{1}{\sqrt{2\pi|C|}} \exp\left\{-\frac{1}{2}n^{\top}C^{-1}n\right\}$$

Assume Gaussian noise*

$$p(\tilde{n}(f)) \sim \frac{1}{\int 2\pi S_n(f) df} e^{-\frac{1}{2} \langle n|n \rangle}$$

Covariance is diagonal (PSD S_n) in frequency domain

$$\langle a|b\rangle = 2\int_0^\infty \frac{\tilde{a}(f)\tilde{b}^*(f) + \tilde{a}^*(f)\tilde{b}(f)}{S_n(f)}df$$

"Noise-weighted inner product"

* noise is not truly Gaussian current detectors

Matched Filtering

$$\hat{d}(t) = \int_{-\infty}^{\infty} d(t') K(t-t') dt'$$

Filter the data with a template *K* that maximizes the signal-to-noise ratio

$$\tilde{K}(f) = c \frac{\tilde{h}(f)}{S_n(f)}$$

Optimal filter is the signal weighted by the PSD

$$ho^2(t) = rac{\langle d | h
angle^2}{\langle h | h
angle}$$

"Matched-filter signal-to-noise ratio"

Matched Filtering in Action



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The first gravitational-wave detection:

GW150914

LIGO Scientific and Virgo Collaborations (2016), Phys. Rev. Lett. 116, 061102

Noise can produce high SNR...

...but loud GW will produce even higher SNR!



LIGO Scientific and Virgo Collaborations (2016), Phys. Rev. Lett. 116, 061102

~ 30 M_o BHs ^{*}/_× ~ 0.5 Gpc away (z~0.1) ^{*}/_×

How do we measure source properties, including uncertainties?



Source Parameters 9: Masses (m_1, m_2) and 3-D dimensionless spin vectors (χ_1, χ_2) of the two coalescing objects, luminosity distance, sky position,...

Data d: Strain in all operating detectors

Likelihood: Gaussian in residuals between strain data and model

Priors: Up to analyst, but we will revisit...

Parameter Estimation in Practice

- GW source posterior is 15+ dimensions!
- Draw samples from the posterior distribution
- Use samples to perform monte carlo integrals over the posterior

 $\int f(\theta) p(\theta|d) d\theta \approx \langle f(\theta) \rangle_{\text{samples}}$



Use **"approximants"** to GR to quickly evaluate h(t) for any source parameters $(m_1, m_2, ...)$

- Effective one-body (EOB) family: include strong-field effects in test particle limit + calibration to NR
- IMRPhenom family: stitch PN and EOB results to NR
- Surrogate: interpolate NR simulations

Catalog of Gravitational Wave Detections



https://git.ligo.org/zoheyr-doctor/plot-gracedb-events

GW170817: Neutron-Star Merger w/ EM counterpart!





Credit: University of Warwick/Mark Garlick



Soares-Santos,..., ZD,... ApJL 848 L16 (2017)

GW170817 Multi-messenger ApJL 848 L12 (2017)

Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars Solar Masses 20 10

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GW200129_065458 GW200202_154313

Parameters for all GW detections...with uncertainties!

GW200216_220804 GW200219_094415 GW200220_061928 GW200220_124850 GW200224_222234 GW200225_060421 GW200302_015811 GW200306_093714 GW200308_173609 GW200311_115853 GW200316_215756 GW200322_091133



LVK, Phys. Rev. X; 13(4):041039
Joint posterior samples for all events and parameters



LVK, Phys. Rev. X; 13(4):041039

GW190521 - Heavy BH Merger!





LVK, ApJL, 900, 1, id.L13, 27 pp.

BINARY BLACK HOLE MERGERS AS OF MAY 2021: 48

FIRST OBSERVATIONS OF BLACK HOLE

NEUTRON STAR MERGERS

BINARY NEUTRON STAR MERGERS AS OF MAY 2021: 2

Credit: Carl Knox, OzGrav/Swinburne University



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@astronerdika Credit: Shanika Galaudage

Updated 2023-10-11	— 01	— 02	— O3	— O4	— O5
LIGO	80 Mpc	100 Мрс	100-140 Мрс	150 160+ Mpc	240-325 Мрс
Virgo		30 Мрс	40-50 Мрс	40-80 Мрс	150-260 Mpc
KAGRA			0.7 Mpc	1-3 ≃10 ≳10 Мрс Мрс Мрс	25-128 Мрс
G2002127-v21 2	I I 2015 2016	I I 2017 2018 2	1 I I 019 2020 2021	I I I I I 2022 2023 2024 2025 2026	I I I 2027 2028 2029



FRANCISCO LLAMAS



AARON JONES

ASHINI MODI



DRIPTA BHATTACHARJEE

Post a Comment





JOCELYN READ



DEBNANDINI MUKHERJEE
Post a Comment



MAYA FISHBACH Post a Comment



humansofligo.blogspot.com

What have we learned from individual events?



- BBH, BNS, NSBH can merge in a Hubble time!
- Some merging black holes spin
- The merging objects can have unequal masses
- BH and NS from GWs have different properties than those observed through EM

How are compact-object mergers produced?

First we need to form compact objects...

Stellar Progenitors of Compact Objects

- Feature at transition from neutron stars to black holes?
- Feature at pair-instability supernova mass?
- Dips/peaks from non-linear mass compactness relation of progenitor stars?



Pair instability supernova (credit NASA)

More Exotic Compact-Object Formation Scenarios

Hierarchical formation

Primordial Formation



Credit: ESA

...then we need compact objects to merge in a Hubble time.

What DOESN'T Work...



The Big Mystery...

Peters (1964):

$$a = \left(\frac{64G^3}{5c^5}M^3 t_{\rm merge}\right)^{\frac{1}{4}} \sim 50 \,\mathrm{R}_{\odot} \,\left(\frac{M}{60 \,\mathrm{M}_{\odot}}\right)^{\frac{3}{4}} \left(\frac{t_{\rm merge}}{14 \,\mathrm{Gyr}}\right)^{\frac{1}{4}}$$

To merge stellar-mass COs in a Hubble time, they must be closer than the radii of their progenitor stars!

Two Families of Compact-Object Merger Channels

Dynamical Channel

Isolated Binary Channel









Belczynski, Holz, Bulik & O'Shaughnessy. Nature (2016)

Many models, many knobs!

- Initial mass function of stars / COs
- Stability of mass transfer
- Cluster potential
- Metallicity evolution
- Accretion efficiency
- ...

Double Compact Object Formation Depends Strongly on Metallicity

More metals... \rightarrow more lines \rightarrow more stellar winds \rightarrow smaller compact objects





High-Dimensional Source Parameter Space!

Noisy Data!

Systematic Uncertainties!

Should explain full catalog of GWs!



Detailed Astrophysical Processes





Detailed Astrophysical Processes





Simple Models

- Targeted questions (e.g. mass gaps)
- Easy to write down (e.g. power laws)
- Somewhat agnostic to astrophysical details
- Could miss important features



Detailed Astrophysical Models

- Can include our best understanding of BHs + interactions
- Can be tuned via other data sets
- Hard to write down
- Many parameters
- Possible systematic errors

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Data-Driven Models

- Find unexpected features in the population
- Corroborate results of simple models
- Compare with features in detailed models
- Black-box predictions

Simple Models Targeted questions Easy to write down Somewhat agnostic to astrophysical details Could miss important features

Strain data segments thatAstrophysical Models +trigger GW search pipelinesSelection Effects

$$\mathcal{L}(\{d\}, N_{\text{det}} | \Lambda, N_{\text{exp}}) \propto$$

$$N^{N_{\text{det}}} e^{-N_{\text{exp}}} \prod_{i=1}^{N_{\text{det}}} \int \frac{\mathcal{L}(d_i | \theta) \pi(\theta | \Lambda) d\theta}{\sum_{\substack{\text{single} \\ \text{event} \\ \text{likelihood}}} \sum_{\substack{\text{single} \\ \text{prior under } \Lambda}}$$

Higher mass mergers are "louder" -> Selection effect



Mass Spectrum of Compact Object Mergers

Structure in the binary black hole mass distribution



LVK PHYS. REV

Spin Spectrum of Compact Object Mergers

Binary black hole spin distribution



LVK PHYS. REV. X 13, 011048 (2023)

Merger rate with redshift

Merger rate is increasing with redshift



LVK PHYS. REV. X 13, 011048 (2023)

Population-level correlations

"Data-Driven" Models

Models with lots of flexibility but "agnostic" to the astrophysics enable...

- Finding unexpected features in the population
- Corroborating results of simple models
- Comparison with features in detailed models
- Black-box predictions for other applications

Let's use these different modeling approaches to study gravitational-wave populations!


Simple Models

Fishbach & Holz (2017) Talbot & Thrane (2018) Wysocki, Lange, and O'Shaughnessy (2019) Doctor et al (2020) Kimball et al (2021) Landry & Read (2021) Farah, Fishbach, Essick et al (2021)

Does the mass distribution of black holes change with redshift?



Fishbach, **ZD**, et al. ApJ 912 98 (2021) 74

Does the mass distribution of black holes change with



Do BH sub-populations of spin have different masses?

Dynamical Channel

Isolated Binary Channel



Do BH sub-populations of spin have different masses?



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Baibhav, ZD, Kalogera (2023)

Do BH sub-populations of spin have different masses?



Simple Models

The Stochastic GW Background from NS Mergers



Bellie, Banagiri, **ZD**, Kalogera. arXiv: 2310.02517 (2023) ⁷⁹

Data-Driven Models

Mandel et al (2016) Farr et al (2018) Powell et al (2019) Tiwari (2021) Rinaldi et al (2021) Sadiq et al (2021) Godfrey et al (2023)

Is there structure in the BBH mass distribution?



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Is there structure in the BBH mass distribution?





Bruce Edelman Research Software Engineer (UOregon)

Data-Driven Models

Edelman, ZD, Godfrey, Farr (2022) ⁸²

Is there structure in the BBH spin distribution?



Detailed Models

BPASS			
BSE			
CMC			
ComBinE			
COMPAS			
COSMIC			
MOBSE			
POSYDON			
SEVN			
StarTrack			

The Isolated Binary Channel



How do merging BH masses compare to those in HMXBs? 175



How do merging BH masses compare to those in HMXBs?

- \lesssim 3% of detectable HMXBs have a BH with > 35 M $_{\odot}$
- Probability detected HMXB will merge as a BBH in a Hubble time is $\lesssim 1\%$
- Discrepant BH masses from GWs and HMXBs are expected!



Minimize how many simulations are needed





Kyle Rocha CIERA grad

locha, ..., ZD, ... (2022)

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A different problem with a similar solution

ZD, Farr, Holz, Puerrer (2017)

Do High-Spin HMXBs become High-Spin Merging BBHs?



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Do High-Spin HMXBs become High-Spin Merging BBHs?



Gallegos-Garcia, Fishbach, Kalogera, Berry, ZD (2022)

Remarks

- Multiple population modeling approaches enable us to understand CO mergers from different angles
- Detailed models: Incorporate our full astrophysical picture, but expensive and many systematics
- Simple parametric models: Empirically test specific questions
- Data-driven models: Look for the unexpected
- Team effort! A rich set of problems for everyone to get involved in
 - Stars, dynamics, statistics, machine learning, detectors and instrumentation...





to confusion, we use merging BBHs for clarity. To identify high-spin HMXBs in simulations, we assume the spin of the first-born BH is imparted by the scenario of Case-A mass transfer (MT) while both stars are on the main sequence (MS; Valsecchi et al. 2010; Qin et al. 2019). In this scenario, the donor star, which is also the progenitor of the first-born BH, could form a highspin BH following a combination of (i) MT that prevents significant radial expansion; (ii) strong tidal synchronization at low orbital periods, and (iii) inefficient AM transport within the massive star post MS. We do not C 11 4.1 1 . DIT

