

Classical Stability and Quantum Effects of Warped AdS₃ Black Holes in Topologically Massive Gravity

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- 2 Classical stability of warped AdS_3 black holes
- 3 QFT on warped AdS_3 black holes
- 4 Conclusions

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Deform theory: **Topologically Massive Gravity**

$$S = S_{\text{E-H}} + S_{\text{C-S}},$$

with:

$$S_{\text{E-H}} = \frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right),$$
$$S_{\text{C-S}} = \frac{\ell}{96\pi G\nu} \int d^3x \sqrt{-g} \epsilon^{\lambda\mu\nu} \Gamma_{\lambda\sigma}^{\rho} \left(\partial_{\mu} \Gamma_{\rho\nu}^{\sigma} + \frac{2}{3} \Gamma_{\mu\tau}^{\sigma} \Gamma_{\nu\rho}^{\tau} \right).$$

Deser, Jackiw, Templeton (1982)

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- **Massive** propagating degree of freedom.
- **Third-order** derivative theory.
- GR solutions \subset TMG solutions.

Warped AdS₃ black hole

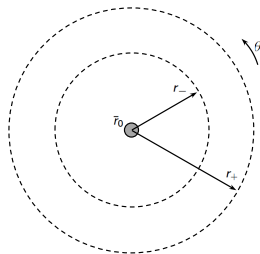
Spacelike stretched black hole:

$$ds^2 = dt^2 + \frac{\ell^2 dr^2}{4R^2(r)N^2(r)} + 2R^2(r)N^\theta(r)dtd\theta + R^2(r)d\theta^2$$

$$R^2(r) = \frac{3(\nu^2 - 1)}{4}r(r - r_0)$$

$$N^2(r) = \frac{(\nu^2 + 3)(r - r_+)(r - r_-)}{4R(r)^2}$$

$$N^\theta(r) = \frac{2\nu r - \sqrt{r_+ r_- (\nu^2 + 3)}}{2R(r)^2}$$

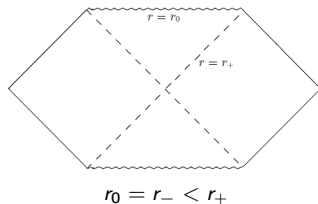
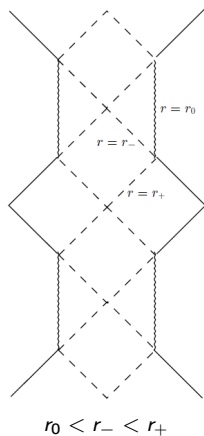


Anninos, Li, Padi, Song, Strominger (2008)

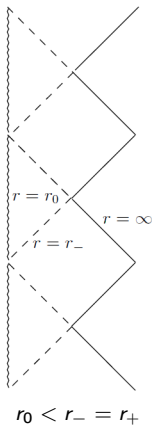
$\nu > 1$ is the **warp factor** of the spacetime.

$\nu \rightarrow 1$ recovers the **BTZ black hole**.

Causal structure of warped AdS₃ black holes



- **Not asymptotically AdS₃!**
- Similar to asymptotically flat black holes!
- Arena to obtain valuable insights for difficult problems with the Kerr black hole!



Jugeau, Moutsopoulos, Ritter (2010)

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- **Scalar field** Φ on the background of a spacelike stretched black hole:

$$(\nabla^2 - m^2) \Phi(t, r, \theta) = 0$$

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- **Exact mode solutions:**

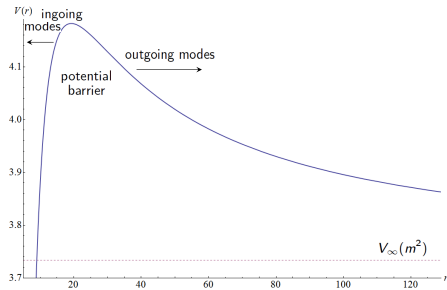
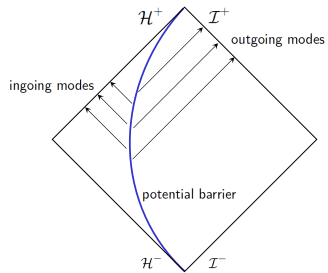
$$\Phi_{\omega k}(t, r, \theta) \sim e^{-i\omega t + ik\theta} z^\alpha (1-z)^\beta F(a, b, c; z)$$

$$z = \frac{r - r_+}{r - r_-}$$

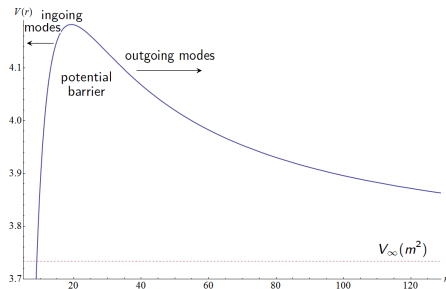
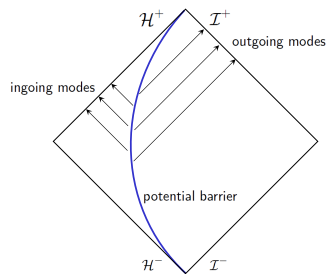
α, β, a, b, c functions of ω and k

$F(a, b, c; z)$ hypergeometric function

Quasinormal modes



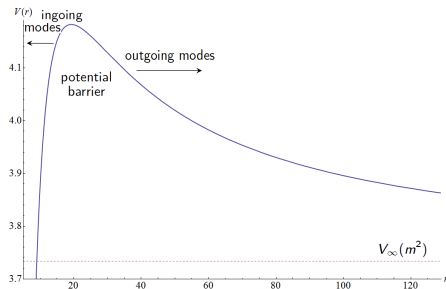
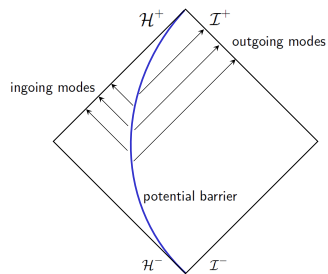
Quasinormal modes



Boundary conditions:

- **Ingoing** modes at the event horizon;
- **Outgoing** modes at infinity.

Quasinormal modes



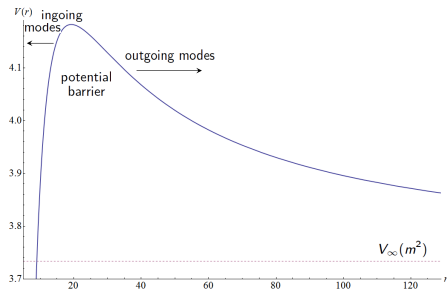
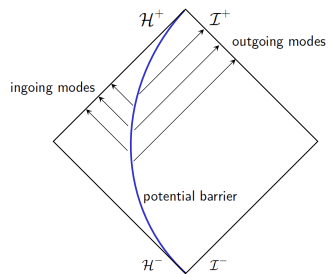
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⇒ Discrete set of complex eigenfrequencies $\{\omega_n\}$

$$\Phi_n \sim e^{-i\omega_n t} = e^{-i\text{Re}(\omega_n)t + \text{Im}(\omega_n)t}$$

Quasinormal modes



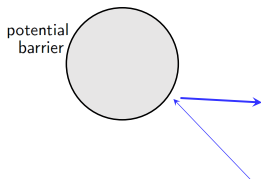
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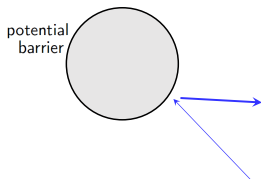
$$\Phi_n \sim e^{-i\omega_n t} = e^{-i\text{Re}(\omega_n)t + \text{Im}(\omega_n)t}$$

If $\text{Im}(\omega_n) > 0$ for some n : mode is **unstable!**



Superradiant modes: amplitude **increases** after reflection by the potential barrier if

$$0 < \text{Re}(\omega) < k\Omega_{\mathcal{H}}$$

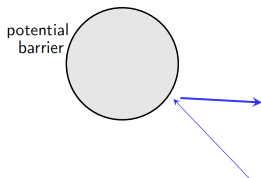


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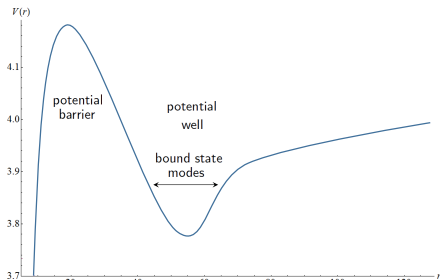
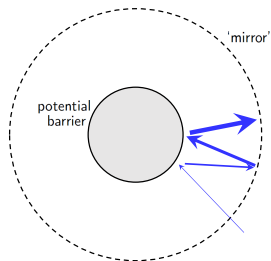
The **existence** of superradiance depends on:

- boundary conditions imposed on the field;
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Whatever the choice of boundary conditions and positive frequency, there are **always** superradiant scalar modes on the warped AdS_3 black hole, similarly to the Kerr black hole (but not to the BTZ and Kerr-AdS).

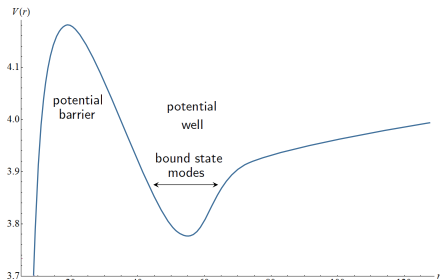
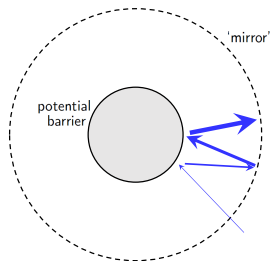
Ferreira (2013)

Superradiant and bound state modes



Bound state modes: localised in the potential well (ingoing at event horizon, exponentially decreasing at infinity).

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$\text{Im}(\omega_n) > 0 \implies$ **superradiant bound state mode** \implies instability!

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NO!

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NO!

- All modes are **stable**: $\text{Im}(\omega_n) < 0$.
- In particular, **no superradiant instabilities**, in contrast with Kerr!

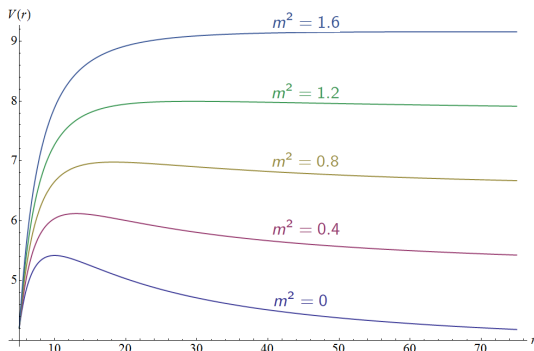
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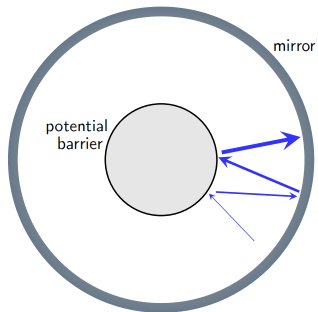
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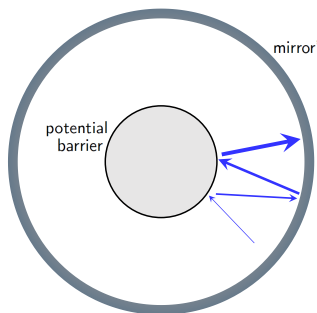
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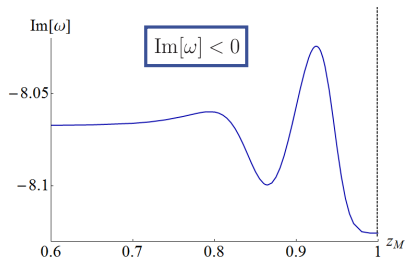
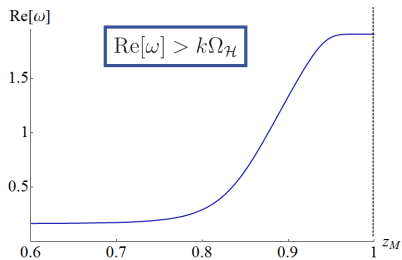


Boundary conditions:

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- **Vanishing** modes at the mirror (Dirichlet boundary condition).

What happens if we add an actual mirror?

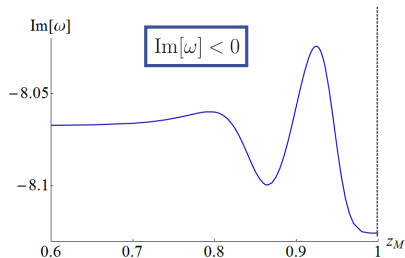
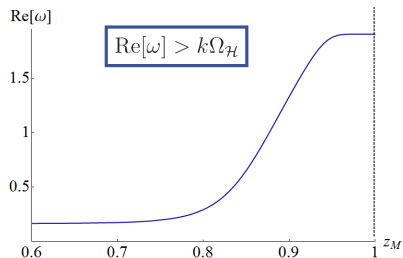
Frequency vs position of the mirror



$$(r_+ = 7, \quad r_- = 0.7, \quad \nu = 1.2, \quad k = -1, \quad m^2 = 0)$$

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Bound state modes are still **stable!**

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- QFT on rotating black holes is a challenging problem:
 - Superradiant modes require care.
 - The Hartle-Hawking vacuum state is **not** well defined!

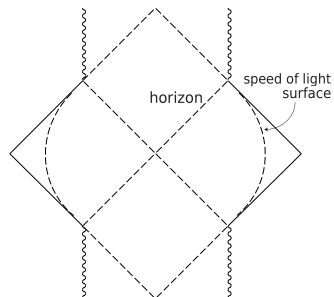
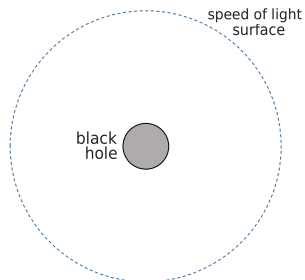
Frolov and Thorne (1989)

Kay and Wald (1991)

Ottewill and Winstanley (2000)

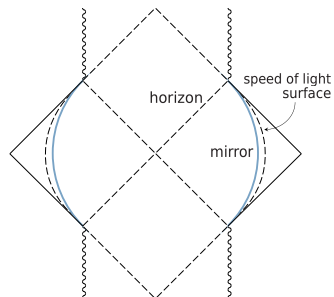
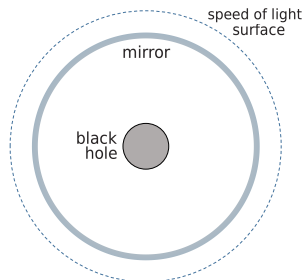
Ottewill and Duffy (2008)

Hartle-Hawking vacuum on a warped AdS_3 black hole



Beyond the speed of light surface, the Hartle-Hawking vacuum would have to rotate with a speed **greater** than the speed of light.

Hartle-Hawking vacuum on a warped AdS_3 black hole



Beyond the speed of light surface, the Hartle-Hawking vacuum would have to rotate with a speed **greater** than the speed of light.

If a **mirror** is put between the horizon and the speed of light surface, an Hartle-Hawking vacuum is **well defined**.

- **Aim:** compute the expectation value of the renormalised stress-energy tensor $\langle T_{\mu\nu}(x) \rangle_{\text{ren}}$ for a scalar field in the Hartle-Hawking vacuum

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Work in progress...

Consider the **complex Riemannian section**:

- Change to **rotating** coords: $(t, r, \theta) \rightarrow (\tilde{t} = t, r = r, \tilde{\theta} = \theta - \Omega_{\mathcal{H}} t)$
- Analytically continue: $\tilde{t} = -i\tau$;
- Impose periodicity: $\tau \sim \tau + \frac{2\pi}{\kappa_+}$ (Hawking temperature $T_H = \frac{\kappa_+}{2\pi}$).

ds_L^2 Lorentzian metric $\rightarrow ds_{\mathbb{C}}^2$ complex Riemannian metric

Stress-energy tensor in the complex Riemannian section

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Green's functions:

$$G^L(x, x') = \int_0^\infty d\tilde{\omega} \sum_{k=-\infty}^{\infty} G_{\tilde{\omega}k}^L(r, r') \rightarrow G^{\mathbb{C}}(x, x') = \sum_{n=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} G_{nk}^{\mathbb{C}}(r, r')$$

Computation is hoped to be easier in the complex Riemannian section.

Frolov (1982)

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- Warped AdS_3 black holes are classically stable to scalar field perturbations, even if a mirror is added to the spacetime, in contrast with Kerr.
- QFT computations on warped AdS_3 black holes may give valuable insights for the Kerr case.

- Is the warped AdS_3 black hole classically stable to other types of perturbations (namely gravitational perturbations)?

What's next?

- Is the warped AdS_3 black hole classically stable to other types of perturbations (namely gravitational perturbations)?
- What is the renormalised stress-energy tensor for a field in the Hartle-Hawking vacuum state? What information does it provide for the Kerr case?

THANK YOU FOR YOUR ATTENTION!

More information in:

[arXiv:1304.6131](https://arxiv.org/abs/1304.6131) (to appear in Phys. Rev. D)

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